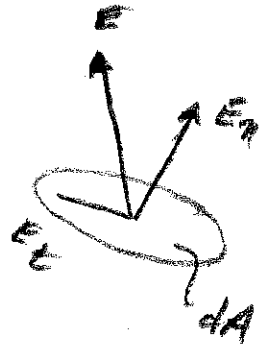


# Gauss' Law

$$\text{Flux} = \Phi = E \cdot A = \int E \cdot \hat{n} \, dA$$



$$\Phi = \int \vec{E} \cdot d\vec{A} = 4\pi k Q_{\text{enc}} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

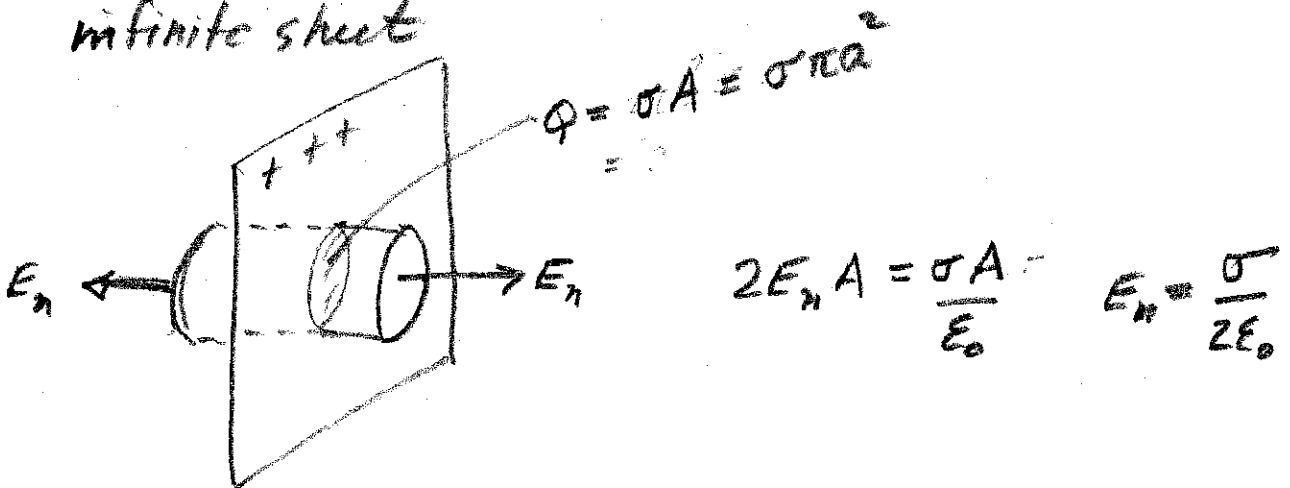
Wed  
03/21/07

Lecture 24

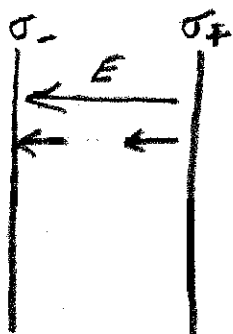
motivation  $E_n = \frac{kQ}{R^2}$

$$\Phi = \frac{kQA}{R^2} = kQ \cdot 4\pi = \frac{Q_{\text{enc}}}{\epsilon_0}$$

infinite sheet

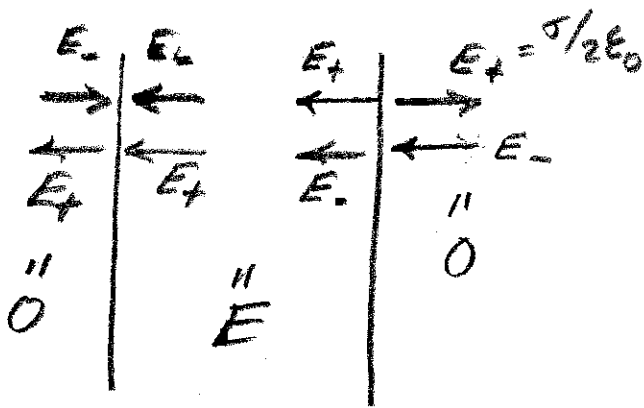


parallel plate capacitor



$$E = \frac{\sigma_+}{2\epsilon_0} + \frac{\sigma_-}{2\epsilon_0} = \frac{\sigma}{\epsilon_0} \text{ inside}$$

$$\int E \cdot dA = 0 \Rightarrow \vec{E} = 0 \text{ outside}$$



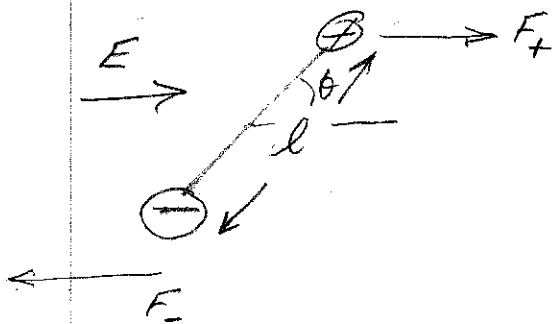
$$F = qE$$

$$F = \sigma E_{dis} = \sigma \left( \frac{\sigma}{2\epsilon_0} \right)$$

$$= \frac{\sigma^2}{2\epsilon_0} \quad \text{don't get fooled}$$

charges can't pull themselves up by their own bootstraps.

### Electric Dipole



$$F_+ = qE$$

$$F_- = -qE$$

$$\Sigma F = 0$$

$$\tau = l \sin\theta F = \sin\theta qlE$$

$$\tau = qlE \sin\theta = pE \sin\theta$$

$$\tau = \vec{p} \times \vec{E} \quad \text{RH-rule}$$

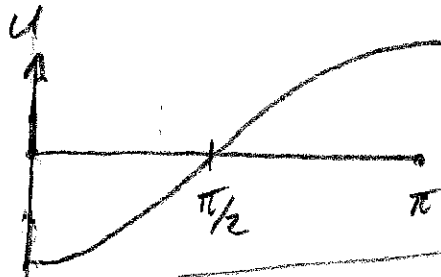
if  $\vec{p}$  defined as  $\ominus \longrightarrow \oplus$

### Potential

$$dW = \tau d\theta = -dU = -pE \sin\theta$$

define  $U=0$  at  $\theta=90^\circ$

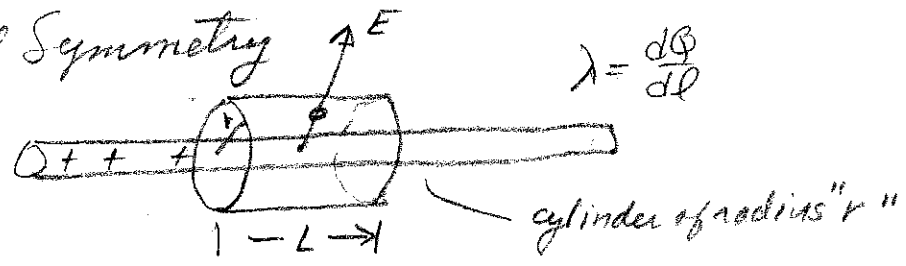
$$U = -pE \cos\theta = -\vec{p} \cdot \vec{E}$$



$$U = qV - \vec{p} \cdot \vec{E} + Q_{ij} \frac{\partial E}{\partial x_j} + \dots$$

Cylindrical Symmetry

wire



$$\lambda = \frac{dQ}{dl}$$

outside

$$\int E \cdot dA = \frac{Q}{\epsilon_0} = E 2\pi r L = \frac{1}{\epsilon_0} (\lambda L)$$

$$E = \frac{\lambda}{2\pi\epsilon_0} \cdot \frac{1}{r}$$

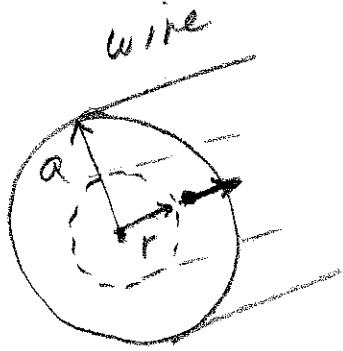
falls as  $1/r$

quad	$\begin{matrix} \oplus & \ominus \\ \ominus & \oplus \end{matrix}$	$= 1/r^4$
dipole	$\begin{matrix} \oplus & \oplus \\ \oplus & \oplus \end{matrix}$	$1/r^3$
mono	$\oplus$	$1/r^2$
line		$1/r$
		const $= Q/\epsilon_0$

Inside

non-conducting  
inside wire

what if  
conducting



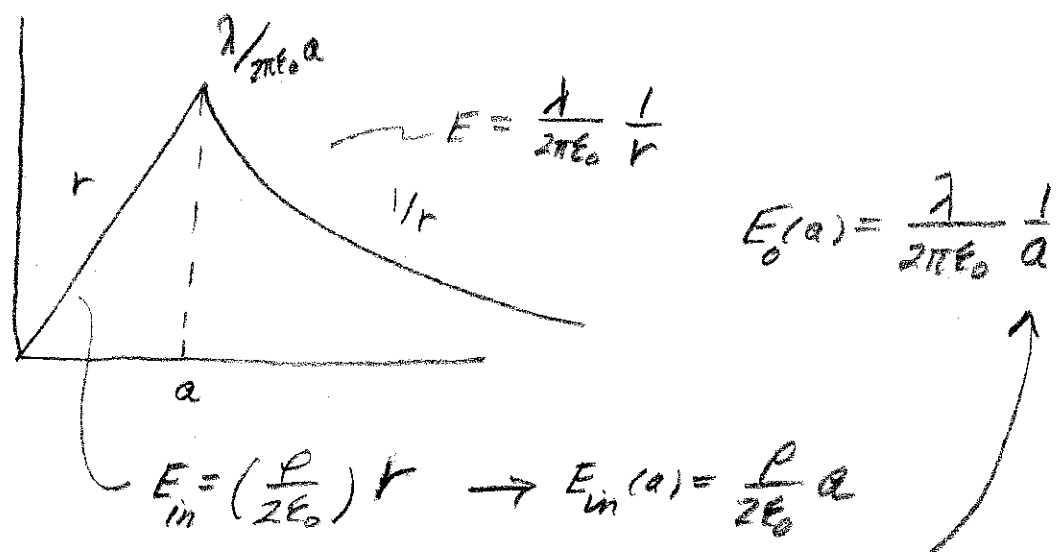
$Q = \lambda L$   
uniform density

$$Q = L \pi a^2 \rho$$

$$\lambda = \frac{Q}{L} = \pi r^2 \rho$$

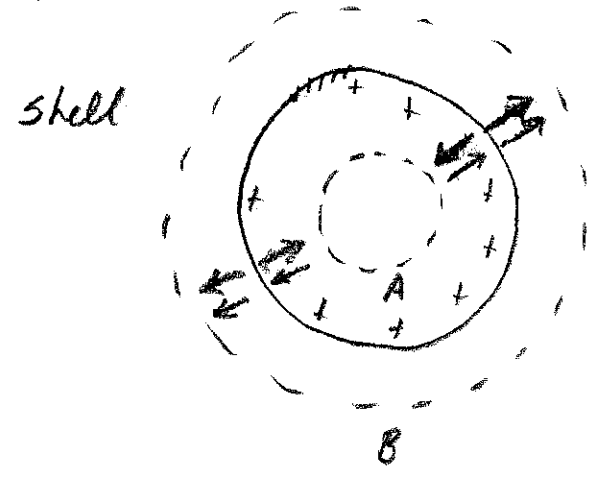
$$E \cdot 2\pi r L = \frac{1}{\epsilon_0} (\pi r^2 \rho L)$$

$$E = \left( \frac{\rho}{2\epsilon_0} \right) \cdot r \quad \text{linear in } r$$



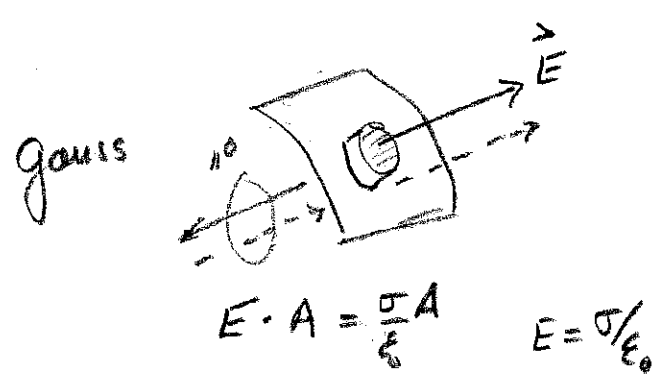
$\rho = \frac{Q}{L\pi a^2} = \frac{\lambda}{\pi a^2}$   
 $E_{in} = \frac{1}{2\epsilon_0} \frac{\lambda}{\pi a^2} \cdot a \Rightarrow \frac{\lambda}{2\pi\epsilon_0 a}$

Spherical Shell - Solid Sphere!

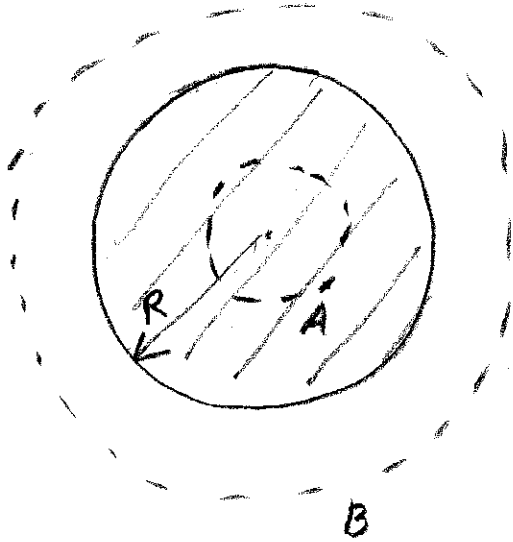


inside = A  
 $E \cdot 4\pi r^2 = \frac{Q}{\epsilon_0} = 0$   
 $\vec{E}_r = 0$

outside = B  
 $E \cdot 4\pi r^2 = \frac{Q}{\epsilon_0} \rightarrow E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$   
 $E = \frac{1}{4\pi\epsilon_0} \frac{4\pi r^2 \sigma}{r^2} = \frac{\sigma}{\epsilon_0}$



Solid Sphere



A = inside

$$E \cdot 4\pi r_A^2 = \frac{1}{\epsilon_0} Q_{\text{enclosed}}$$

$$= \frac{1}{\epsilon_0} \left\{ \rho \frac{4}{3} \pi r_A^3 \right\}$$

$$E = \frac{4\pi \rho r_A^3}{3 \cdot 4\pi \epsilon_0 r_A^2} = \frac{\rho}{3\epsilon_0} r_A$$

$$Q_{\text{tot}} = \frac{4}{3} \rho \pi R^3$$

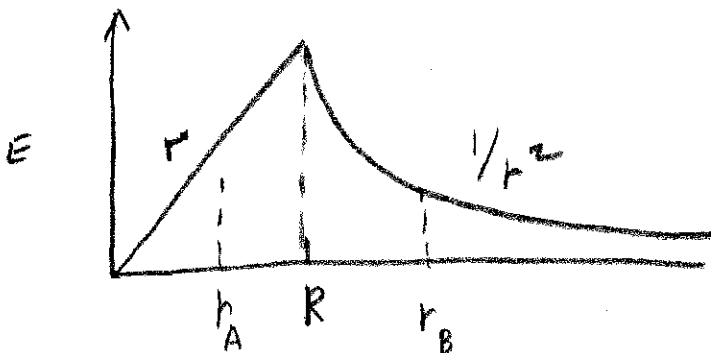
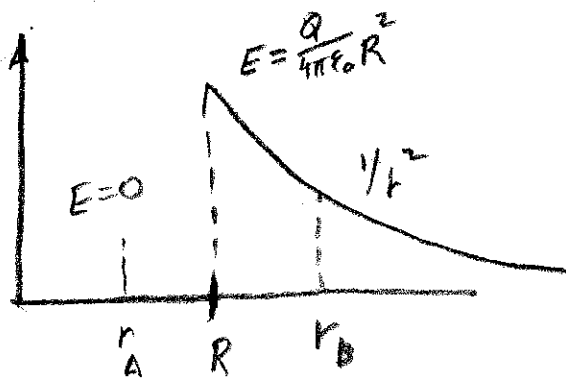
$$E = \frac{Q}{4\pi \epsilon_0 R^3} \cdot r_A \rightarrow \frac{Q}{4\pi \epsilon_0 R^2}$$

at surface

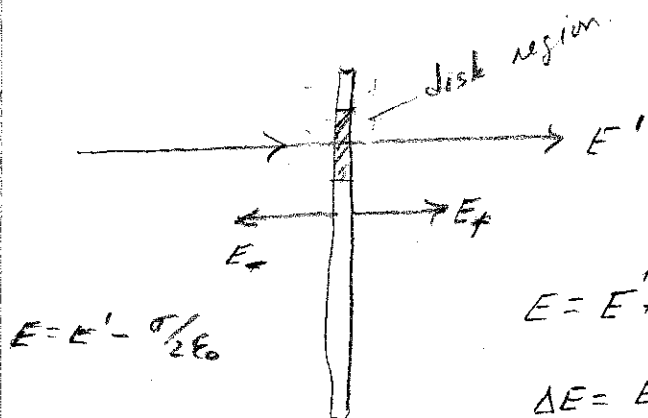
B = outside

$$E (4\pi r_B^2) = \frac{Q}{\epsilon_0}$$

$$E_B = \frac{1}{4\pi \epsilon_0} \frac{Q}{R^2}$$

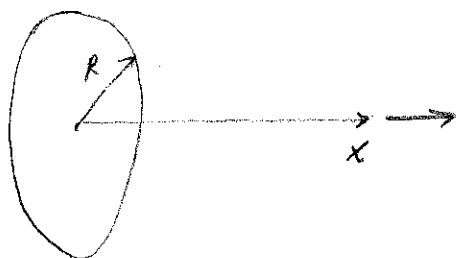


# Discontinuities of $\vec{E}$



$$E = E' + E_t = E' + \frac{\sigma}{\epsilon_0}$$

$$\Delta E = E' + \frac{\sigma}{2\epsilon_0} - (E' - \frac{\sigma}{2\epsilon_0}) = \frac{\sigma}{\epsilon_0} \text{ discontinuity}$$

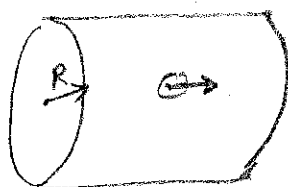


$$E_x = \frac{\sigma}{2\epsilon_0} \left( 1 - \frac{1}{\sqrt{1 + \frac{R^2}{x^2}}} \right)$$

as  $x \rightarrow 0$

$$E_x \rightarrow \frac{\sigma}{2\epsilon_0}$$

cylinder



$$E = \frac{\lambda}{2\pi\epsilon_0 R} = \frac{2\pi R \sigma}{2\pi\epsilon_0 R} \rightarrow \frac{\sigma}{\epsilon_0}$$

Conductors - charges reside on surfaces. (statics)

