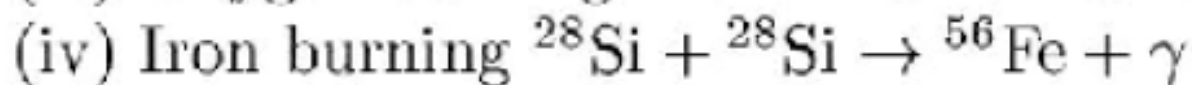
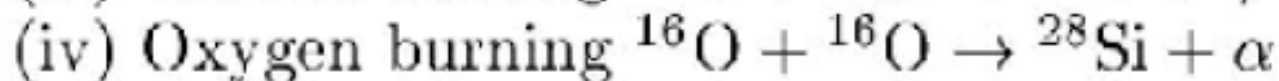
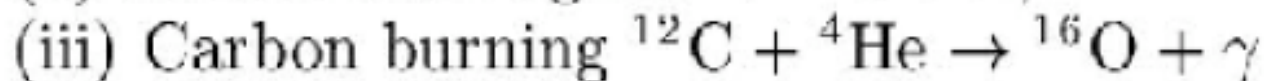
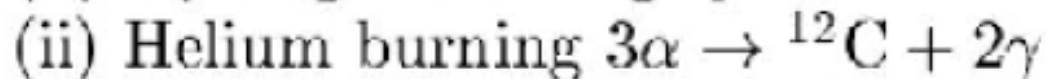
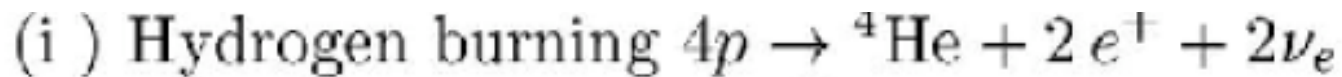


Core Collapse

A star passes most of its lifetime burning H (main sequence). The resulting He builds up in the core and its mass increase, heating and contracting under the pressure of outer layer. The star contraction pauses as nuclear fusion provides the energy necessary to replenish the energy the star loses in radiation and neutrinos. When the T in the core is sufficiently large, He burning begins. After He burning the evolution is greatly accelerated by neutrino losses. The scheme repeats for different stages



Million yrs

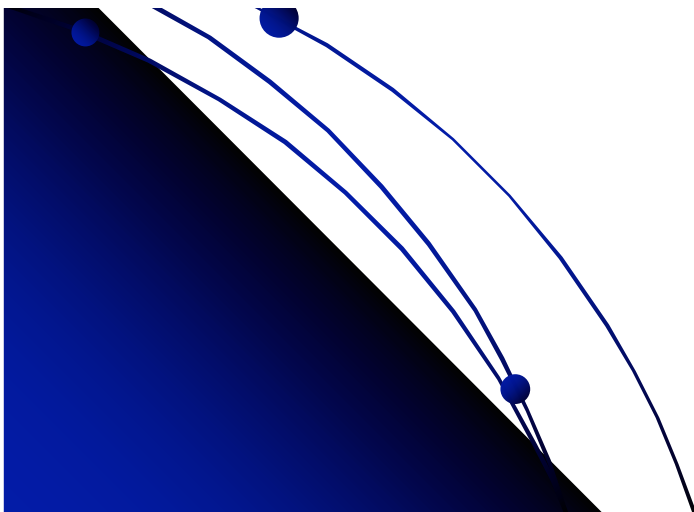
Few weeks



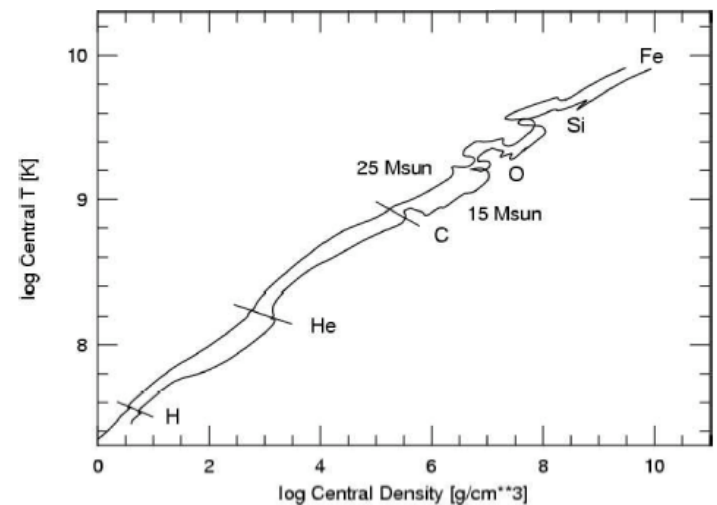
Evolution of a $15 M_{\text{sun}}$ star

Stage	Time Scale	Fuel or Product	Ash or product	Temperature (10^9 K)	Density (gm/cm^3)	Luminosity (solar units)	Neutrino Losses (solar units)
Hydrogen	11 My	H	He	0.035	5.8	28,000	1800
Helium	2.0 My	He	C,O	0.18	1390	44,000	1900
Carbon	2000 y	C	Ne,Mg	0.81	2.8×10^5	72,000	3.7×10^5
Neon	0.7 y	Ne	O,Mg	1.6	1.2×10^7	75,000	1.4×10^8
Oxygen	2.6 y	O,Mg	Si,S,Ar, Ca	1.9	8.8×10^6	75,000	9.1×10^8
Silicon	18 d	Si,S,Ar, Ca	Fe,Ni, Cr,Ti,...	3.3	4.8×10^7	75,000	1.3×10^{11}
Iron core collapse ^a	~1 s	Fe,Ni, Cr, Ti,...	Neutron Star	> 7.1	$> 7.3 \times 10^9$	75,000	$> 3.6 \times 10^{15}$

^aThe presupernova star is defined by the time when the contraction speed anywhere in the iron core reaches $1,000 \text{ km s}^{-1}$.



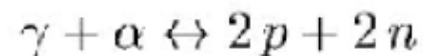
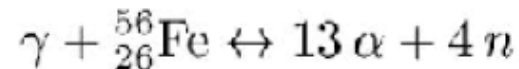
Teresa Montaruli, Apr. 2006



Collapse

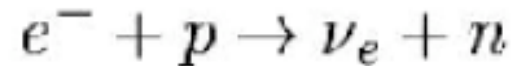
Each cycle requires a higher T for ignition due to the stronger Coulomb repulsion between nuclei. For most stars the process stops when the pressure is not sufficient to heat the core at the necessary T for the next ignition and the star turns into a white dwarf.

The most massive stars can develop an iron core of (iron is the “ground state” of nuclear matter, the most tightly bound of all nuclei and no further nuclear burning is possible). The Fe core is sustained by the degenerate pressure of its electrons until it reaches the Chandrasekar mass of $1.4 M_{\text{Sun}}$. After this limit the core collapses. photodisintegration (radiation melt down some of the Fe nuclei to He) contribute in reducing pressure:



Collapse

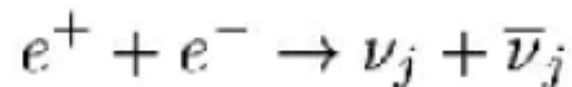
Neutronization follows (time scale of ms): electron capture



About 10^{57} ν_e are emitted contributing to the collapse.

The core collapses to a hot n rich sphere of about 30 km in radius (proto-neutron star).

The **short range nuclear force halts the collapse** when the density is about 2 x atomic nucleus density $4-5 \cdot 10^{14}$ g/cm³. Neutrinos remain trapped in the collapsing core and are in thermal equilibrium within the core. The energy is released in thermal processes like **(thermalization)**



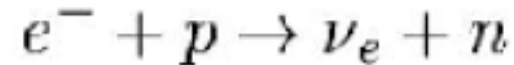
With an **emission of the order of 10 s.**

The shock wave produced by the abrupt halt of the collapse and the bounce of the core travels towards the surface of the star. This is the explosion visible in the optical.



Collapse

Neutronization follows (time scale of ms): electron capture



About 10^{57} ν_e are emitted contributing to the collapse

Neutronization burst of ν_e

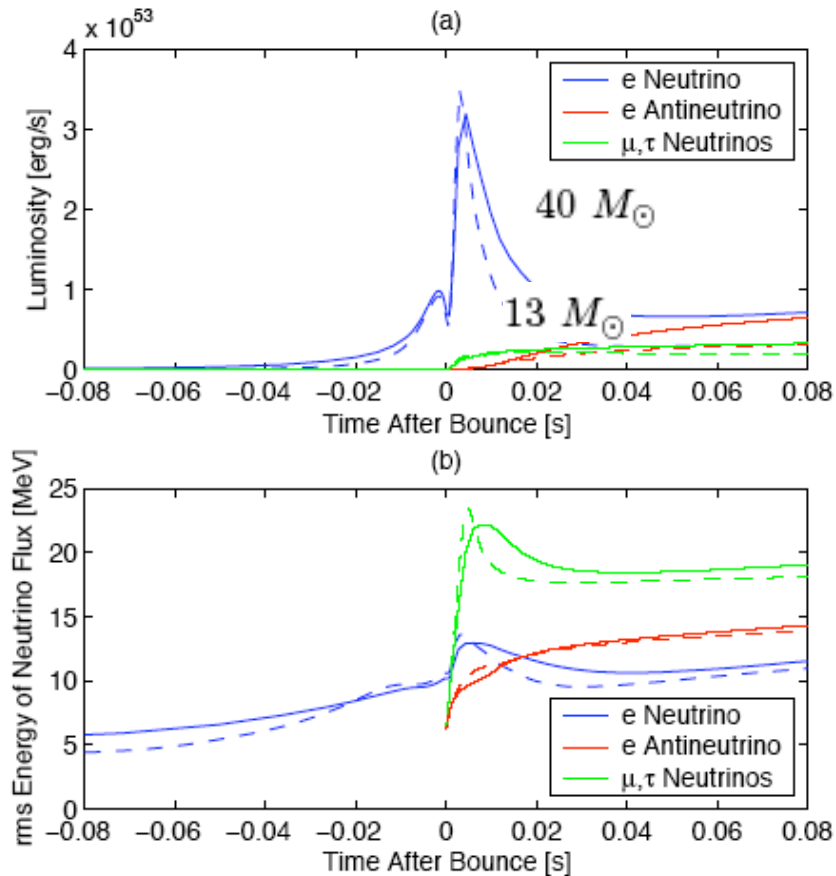
$$N_{\nu_e} \simeq N_e = N_p \simeq \frac{M_{\text{core}}}{m_p} Y_e \simeq 0.9 \times 10^{57}$$

$Y_e \equiv n_{e^-} - n_{e^+}$ is electron fraction per baryon

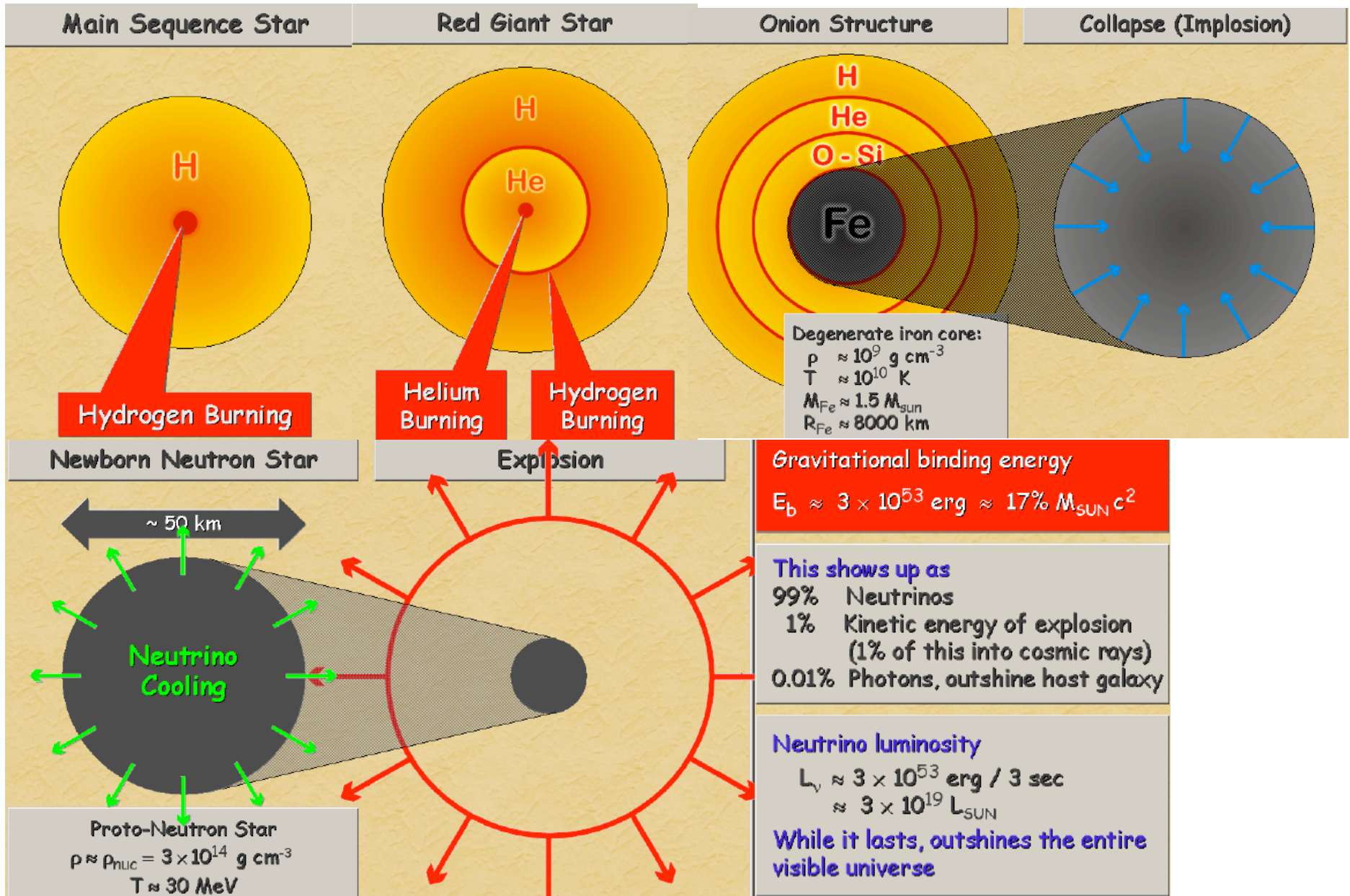
The total energy

$$E_{\text{binding}} = G \frac{M_{\text{core}}^2}{R_f} - G \frac{M_{\text{core}}^2}{R_i} \approx G \frac{M_{\text{core}}^2}{R_f} \approx 3 \times 10^{53} \text{ erg}$$

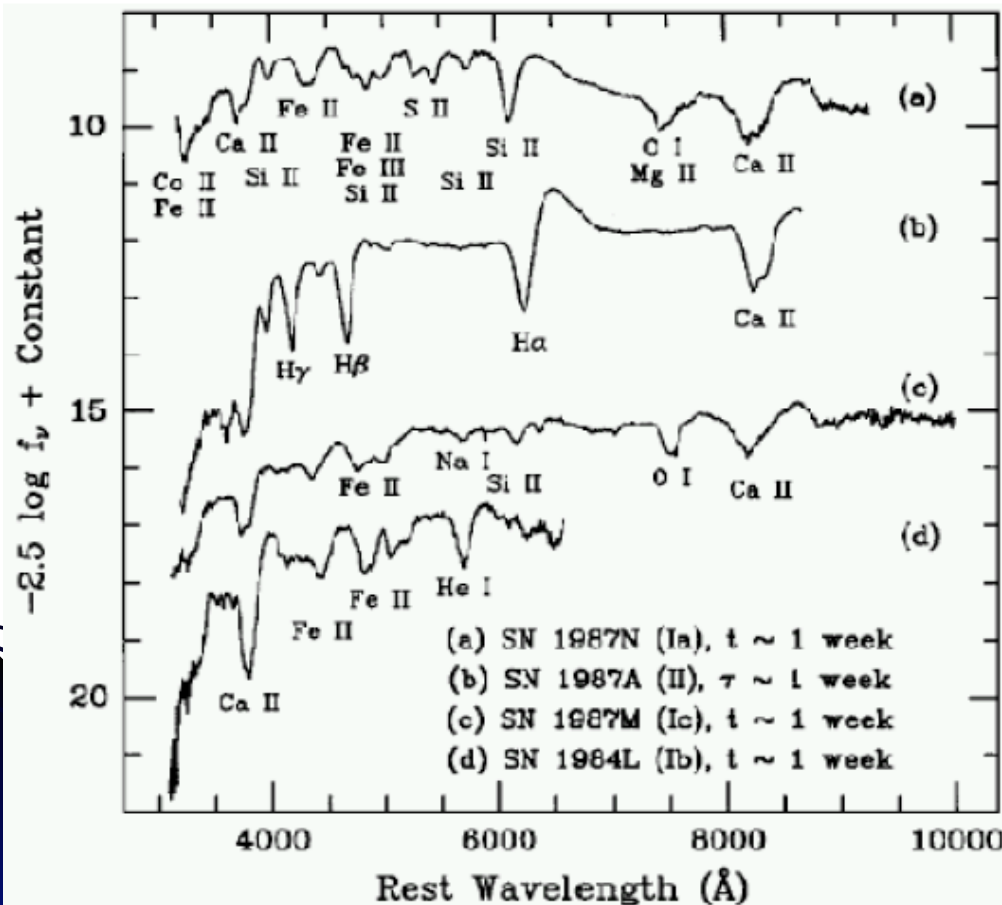
is accounted for by neutrinos while the ejecta carry only 10^{51} erg (1%)



Stellar collapse and SN explosion



Classification of SuperNovae



Type I
(no H)

Type Ia
(no H, strong Si)

Type Ib
(no H, obvious He)

Type Ic
(no H, He, Si)

Type II
(obvious H)

Type II, outer H layer remains at collapse;

Type Ib, outer H layer stripped before collapse;

Type Ic, outer H and He layers stripped before collapse.

Classification of Supernovae

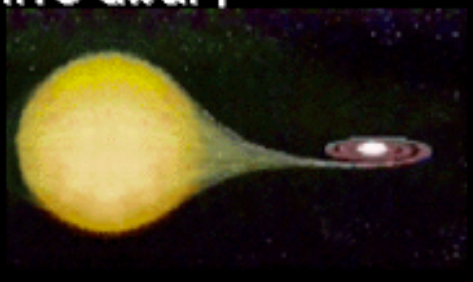
Spectral Type	Ia	Ib	Ic	II
Spectrum	No Hydrogen			Hydrogen
	Silicon	No Silicon		
		Helium	No Helium	
Physical Mechanism	Nuclear explosion of low-mass star	Core collapse of evolved massive star (may have lost its hydrogen or even helium envelope during red-giant evolution)		
Light Curve	Reproducible	Large variations		
Neutrinos	Insignificant	~ 100 × Visible energy		
Compact Remnant	None	Neutron star (typically appears as pulsar) Sometimes black hole ?		
Rate / h ² SNU	0.36 ± 0.11	0.14 ± 0.07		0.71 ± 0.34
Observed	Rate approx. 1 SN / 30 years / galaxy			

$$1 \text{ Snu} = 1 \text{ SN} / 10^{10} L_{\text{sun,B}} / 100 \text{ yrs}$$

Type Ia vs. Core-Collapse Supernovae

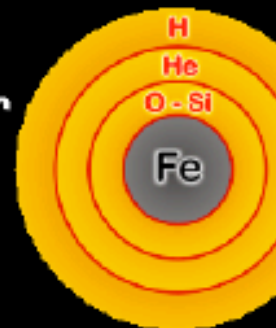
Type Ia

Carbon-oxygen white dwarf
(remnant of
low-mass star)
accretes matter
from companion



Core collapse (Type II, Ib/c)

Degenerate iron core
of evolved massive star
Accretes matter
by nuclear burning
at its surface



Chandrasekhar limit is reached - $M_{Ch} \approx 1.5 M_{sun} (2Y_e)^2$

COLLAPSE SETS IN

Nuclear burning of C and O ignites
→ Nuclear deflagration
("Fusion bomb" triggered by collapse)

Collapse to nuclear density
Bounce & shock
Implosion → Explosion

Powered by nuclear binding energy

Powered by gravity

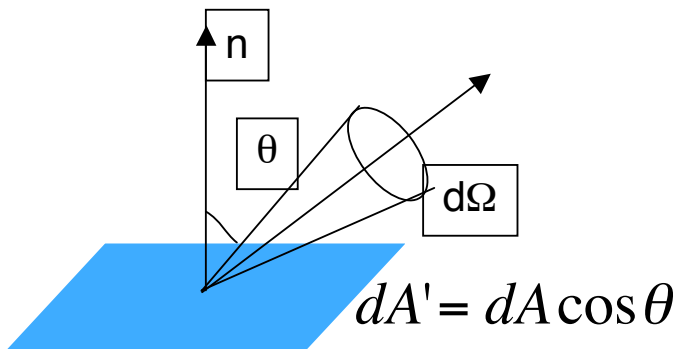
Gain of nuclear binding energy
~ 1 MeV per nucleon

Gain of gravitational binding energy
~ 100 MeV per nucleon
99% into neutrinos

Comparable "visible" energy release of $\sim 3 \times 10^{51}$ erg

Some definitions

Definition of Specific intensity
 Electromagnetic energy dE passing
 Through surface dA and coming
 from an angle θ within a solid angle
 $d\Omega$ during time dt :



$$dE = I_\nu(\theta, \phi) \cos \theta d\nu dA dt d\Omega$$

$$I_\nu = I_\nu(\theta, \phi)$$

$$[I_\nu] = W \cdot m^{-2} sr^{-1} Hz^{-1}$$

The Intensity (integral)

$$I(Wm^{-2}sr^{-1}) = \int_{\nu_1}^{\nu_2} I_\nu d\nu$$

Flux density = intensity integrated over the
 solid angle (of the source) is called:

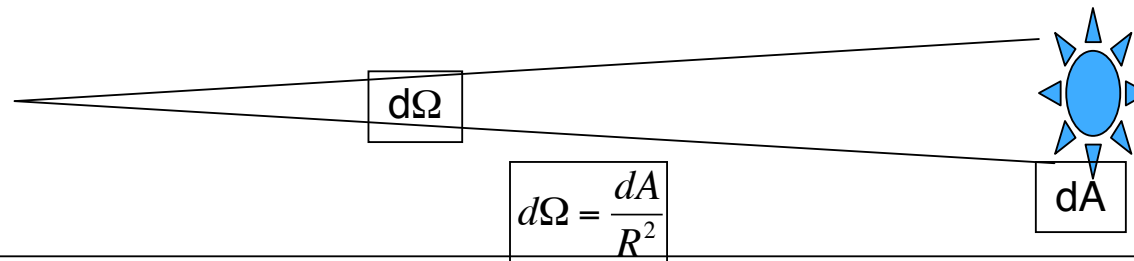
$$dF_\nu = I_\nu \cos \theta d\Omega$$

$$F_\nu = \int_{\Omega} I_\nu \cos \theta d\Omega$$

$$[F_\nu] = W \cdot m^{-2} Hz^{-1}$$

$$1 \text{ Jansky} = 1 \text{ Jy} = 10^{-26} W \cdot m^{-2} Hz^{-1}$$

Flux and Luminosity



- Flux density at telescope

$$F_{\nu} = \int_{\Omega_{Source}} I_{\nu} \cos\theta d\Omega \approx \int_{\Omega_{Source}} I_{\nu} d\Omega$$
$$F_{\nu} = \frac{1}{R^2} \int_{\Omega_{Source}} I_{\nu} dA$$

- Luminosity (isotropic source):

- differential

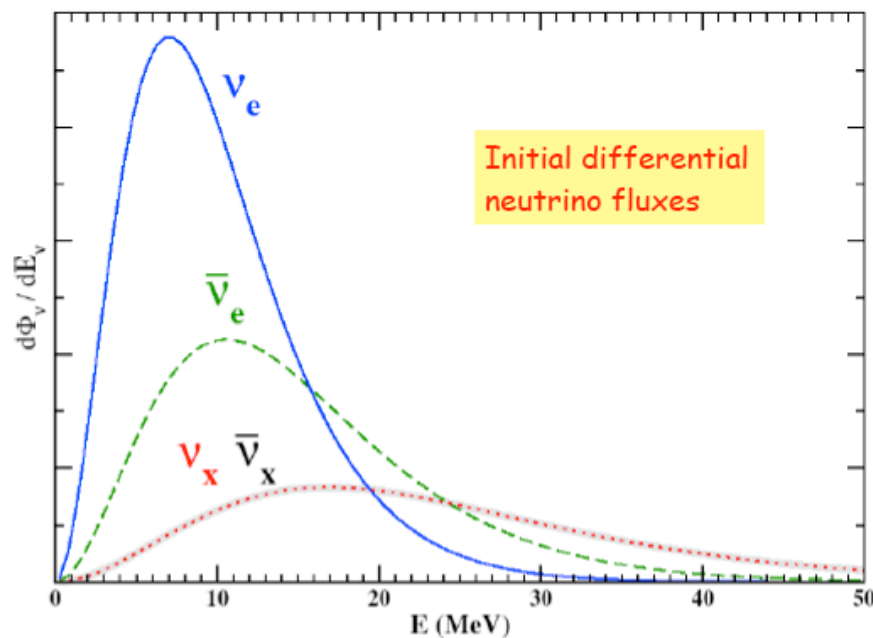
$$L_{\nu} = 4\pi R^2 F_{\nu}$$

- total

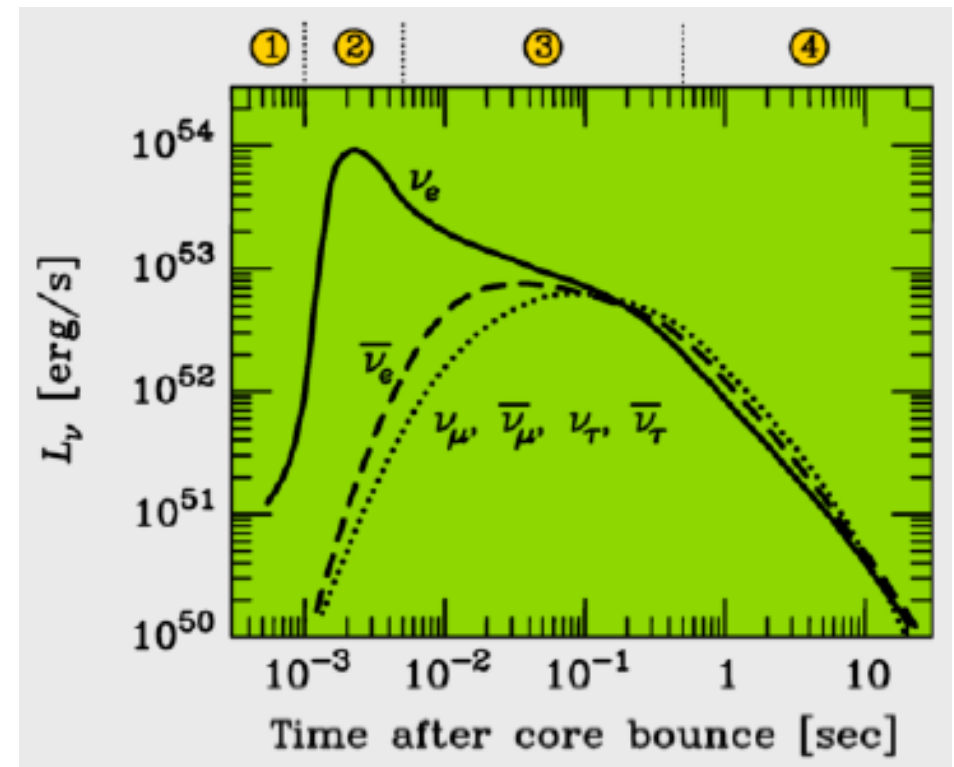
$$L = 4\pi R^2 \int_0^{\infty} F_{\nu} d\nu$$

1. Collapse (infall phase)
2. Shock break out
3. Matter accretion
4. Kelvin-Helmholtz cooling

$$\phi_{\nu_\alpha} \simeq C \frac{E^2}{e^{E/T_{\nu_\alpha}} + 1}$$



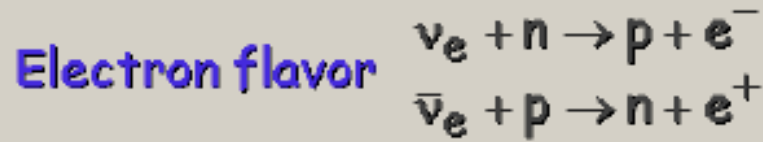
Thermal Spectra



$$E_{\nu_e}^{\text{tot}} = E_{\bar{\nu}_e}^{\text{tot}} = E_{\nu_\mu}^{\text{tot}} = E_{\bar{\nu}_\mu}^{\text{tot}} = E_{\nu_\tau}^{\text{tot}} = E_{\bar{\nu}_\tau}^{\text{tot}} \simeq \frac{E_{\text{binding}}}{6}$$

What determines the time scale?

Main neutrino reactions



Mean free path

$$\lambda = (\sigma n_B)^{-1} \approx 28 \text{ cm} \left(\frac{100 \text{ MeV}}{E_\nu} \right)^2$$

Nucleon density

$$n_B = \frac{\rho_{\text{nuc}}}{m_N} \approx 1.8 \times 10^{38} \text{ cm}^{-3}$$

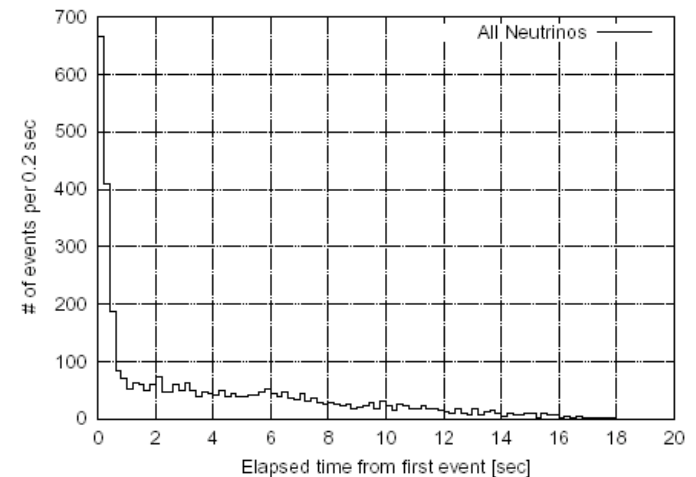
Diffusion time

$$t_{\text{diff}} \approx \frac{R^2}{\lambda} \approx 1.2 \text{ sec} \left(\frac{R}{10 \text{ km}} \right)^2 \left(\frac{E_\nu}{100 \text{ MeV}} \right)^2$$

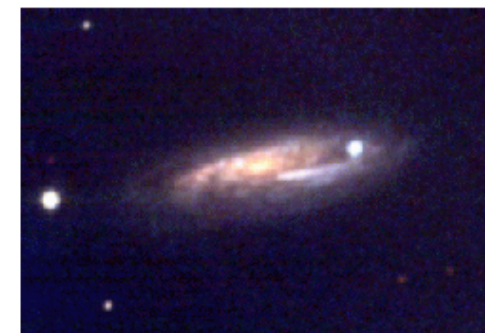
ν_μ, ν_τ only NC so leave with highest $T=8 \text{ MeV}$ and $\langle E \rangle \sim 25 \text{ MeV}$

$\nu_e, \bar{\nu}_e$ also CC hence leave with lower $T=3.5 \text{ MeV}$ and $5 \text{ MeV} \Rightarrow 16 \text{ MeV}$ and 11 MeV . ν_e have lower energy since the material is n rich and thus they interact more

teresa montaruli, Apr. 2000



SN almost as bright as Galaxies!



SN 1998dh
in NGC7541

: 2006



SN 1994D in NGC 4526



SN 1998S in NGC 3877

The historical SNs

Over the past 2000 yrs we have historical records of AD 185, 1006, 1054 Crab Nebula, 1181, 1572 Tycho's SN, 1604 Kepler's SN

<http://arxiv.org/pdf/astro-ph/0301603>

Table 1. Summary of the historical supernovae, and the source of their records

date	length of visibility	remnant	Historical Records				
			Chinese	Japanese	Korean	Arabic	European
AD1604	12 months	G4.5+6.8	few	–	many	–	many
AD1572	18 months	G120.1+2.1	few	–	two	–	many
AD1181	6 months	3C58	few	few	–	–	–
AD1054	21 months	Crab Nebula	many	few	–	one	–
AD1006	3 years	SNR327.6+14.6	many	many	–	few	two
AD393	8 months	–	one	–	–	–	–
AD386?	3 months	–	one	–	–	–	–
AD369?	5 months	–	one	–	–	–	–
AD185	8 or 20 months	–	one	–	–	–	–

八年六月己巳客星出奎宿犯
 傳舍占客星亦妖星天之使者見於天而無常所入列
 宮以示休咎星大者事大而禍深色白其分有兵喪今
 客星出紫微外座傳舍星宜備竊使邊夷侵境又云出
 奎宿為兵竊臣偽惑天子於是金虜遣使來爭執地書
 儀甲戌客星守傳舍第五星 九年正月癸酉客星始
 不見自去年六月己巳至是凡一百八十五日乃消伏
 時虜使久在館至是乃去

Supernova 1054 Petrograph

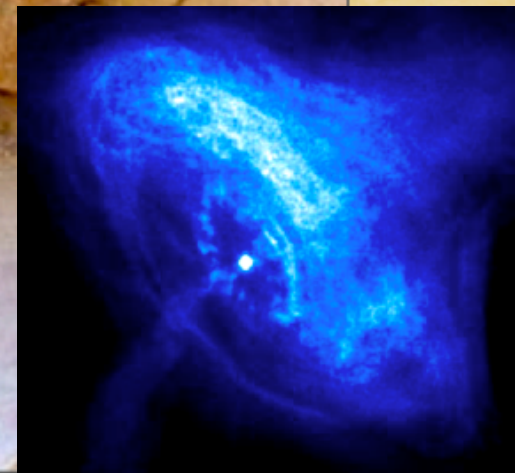
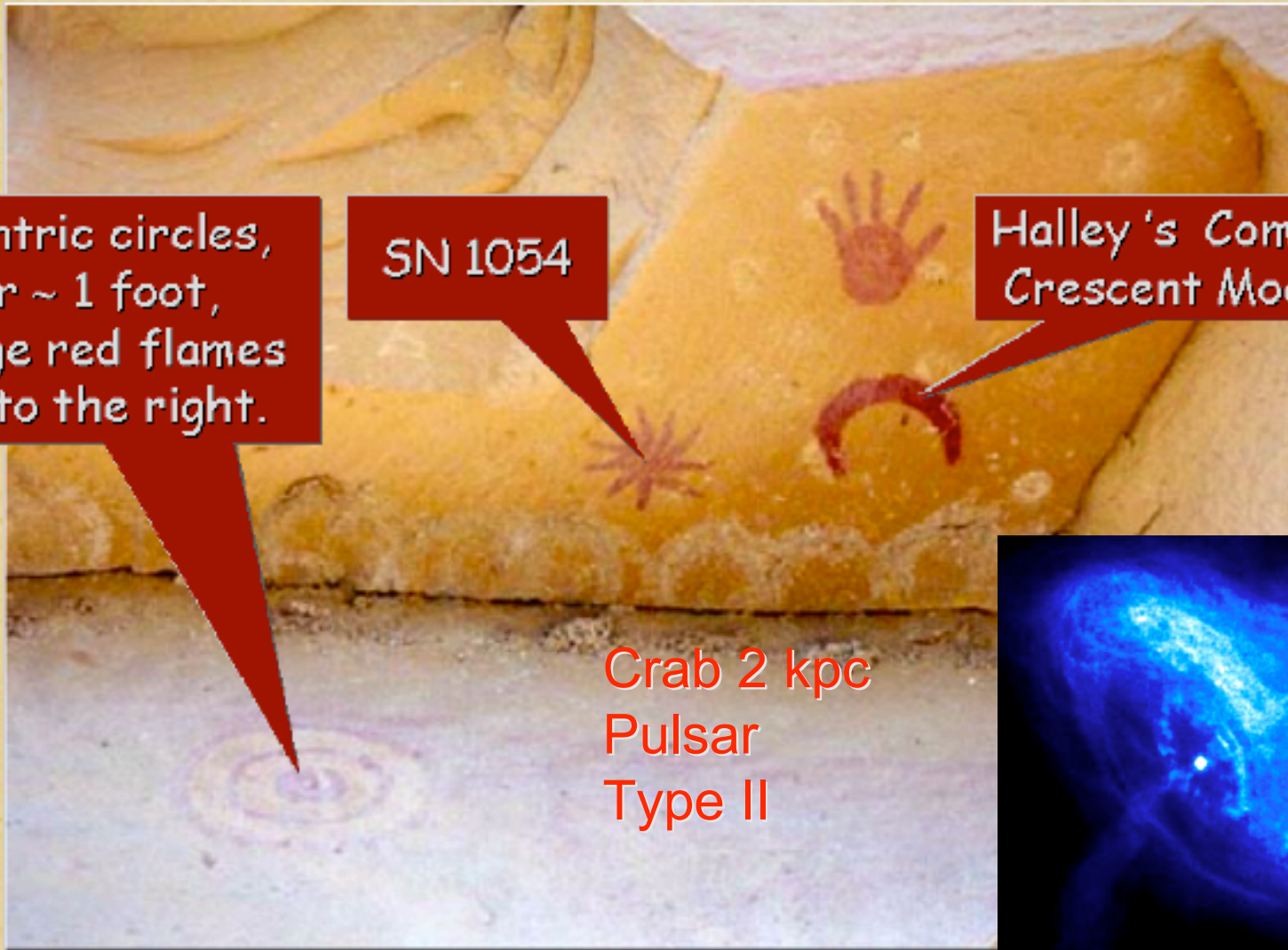
3 concentric circles,
diameter ~ 1 foot,
with huge red flames
trailing to the right.

SN 1054

Halley's Comet?
Crescent Moon?

Crab 2 kpc
Pulsar
Type II

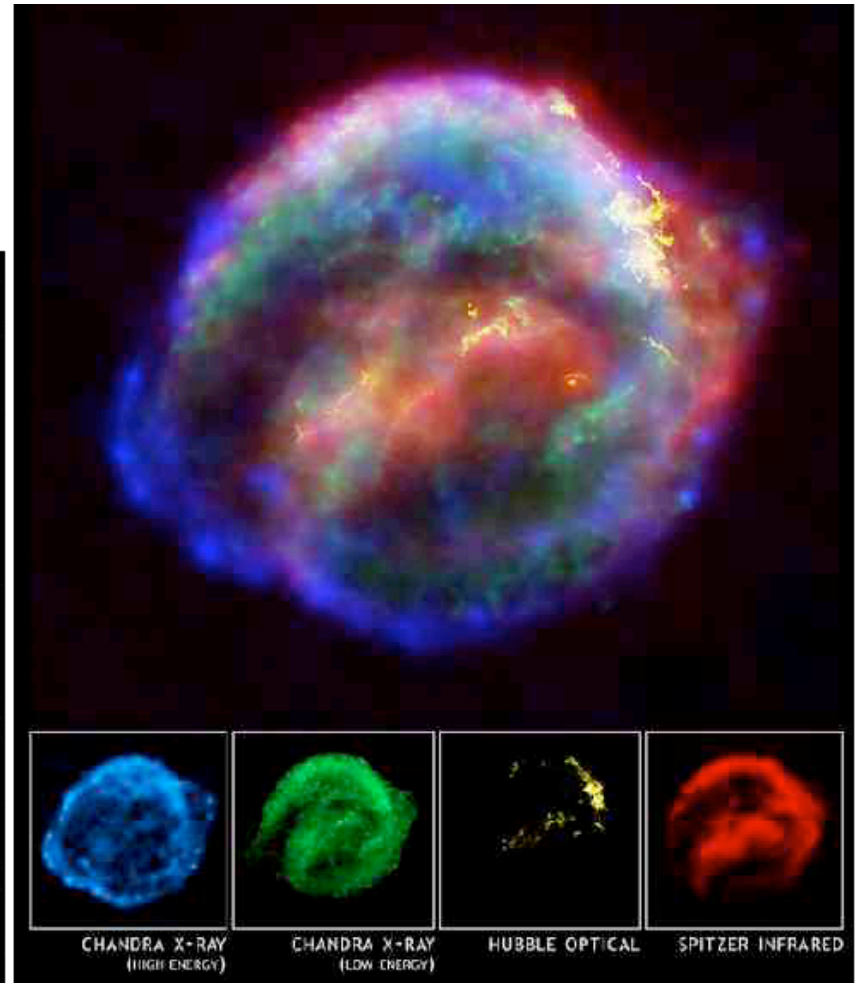
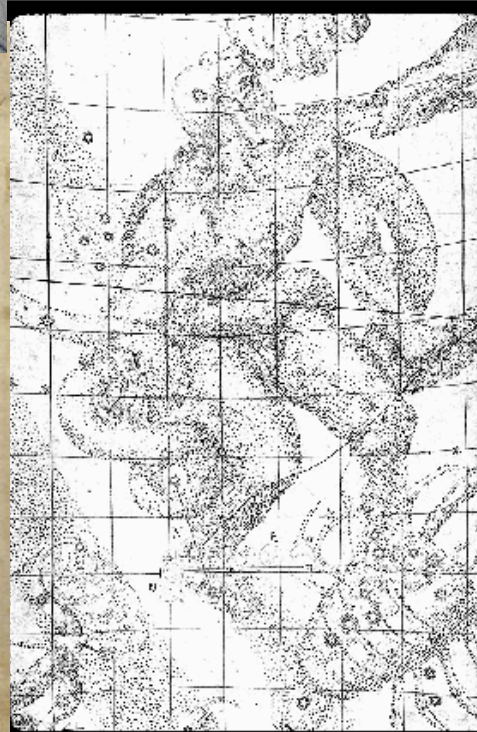
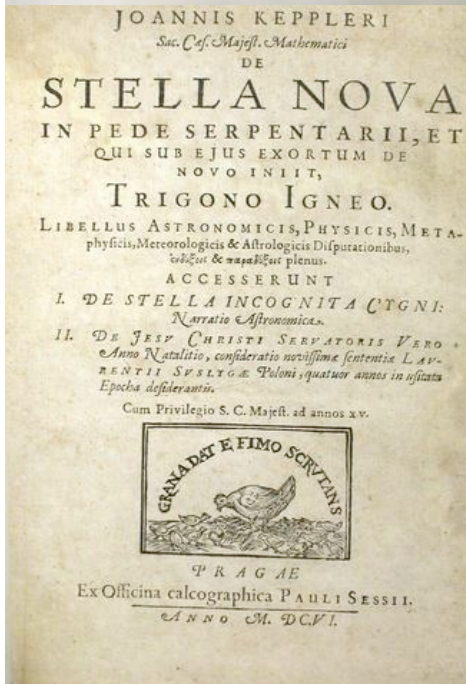
Possible SN 1054 Petrograph by the Anasazi people
(Chaco Canyon, South-Western U.S.)



SN 1604

4-8 kpc Type II?

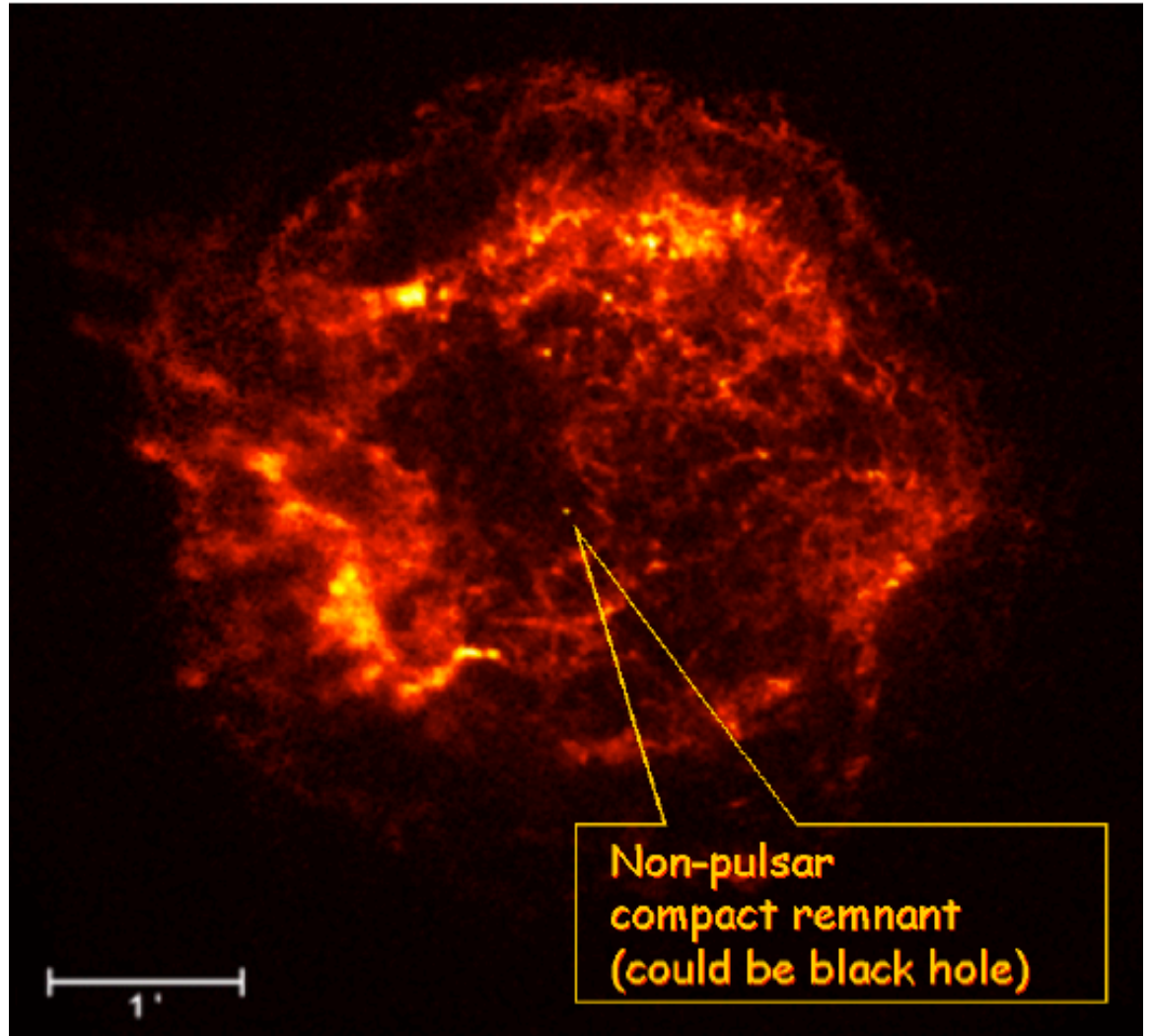
Johannes Kepler,
*De Stella Nova in Pede
Serpentarii*, (1606)



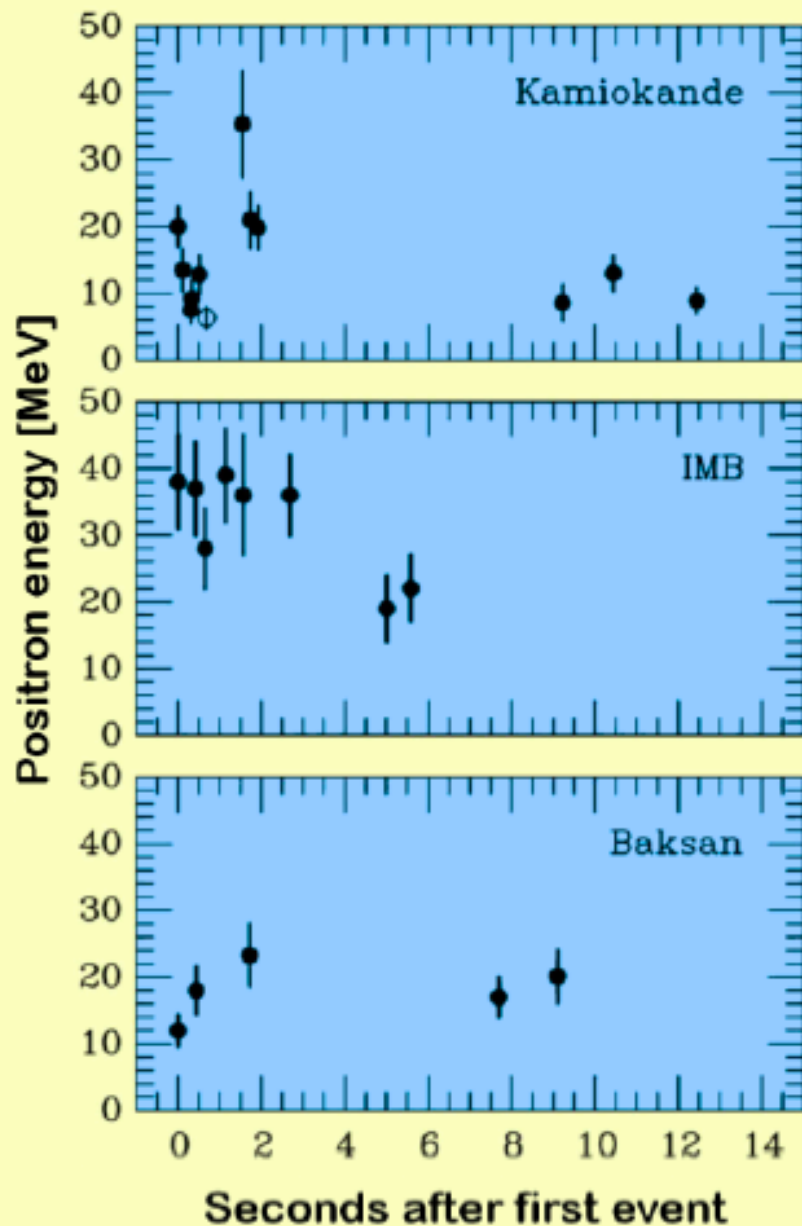
Cassiopeia A

2.8 kpc, neutron star

John Flamsteed
August 16, 1680



Neutrino Signal of Supernova 1987A



Kamiokande (Japan)
Water Cherenkov detector
Clock uncertainty ± 1 min

Irvine-Michigan-Brookhaven (US)
Water Cherenkov detector
Clock uncertainty ± 50 ms

Baksan Scintillator Telescope
(Soviet Union)
Clock uncertainty $+2/-54$ s

Within clock uncertainties,
signals are contemporaneous

Detection of SN neutrinos

$$E_e = E_\nu - 1.3 \text{ MeV}$$

Largest cross section $\bar{\nu}_e + p \rightarrow n + e^+$

~300 e⁺/kt in water

$$\sigma_{\bar{\nu}_e p} \simeq \frac{G_F^2}{\pi} (1 + 3g_a^2) E_\nu^2 \simeq 9.77 \times 10^{-44} \left(\frac{E_\nu}{1 \text{ MeV}} \right) \text{ cm}^2$$

expected rate in Galaxy: 2-4 /century

H₂O Detectors:

SK 22.5 kt (fiducial) 31 Apr 96-15 Jul 01 : search for ν bursts 1704 d $E_{\text{th}} = 6.5$ MeV Expected: 3500 events for 10 kpc SN 12 Msun (2% decrease due to E_{th} changement in SK-II due to 1/2 PMTs), limit on number of explosions/yr: 0.49 SN/yr (90%c.l.) full efficiency up to 100 kpc in SNEWS

AMANDA II 677 Oms $V_{\text{eff}}/OM \sim 400\text{-}500 \text{ m}^3$ 4.3 SN/yr (90%c.l.) in Galaxy (Ahrens et al, 2002): expect 15 fake/yr \Rightarrow SNEWS bckg < 1/week)

SNO 1+1.4 kt (D₂O+H₂O)

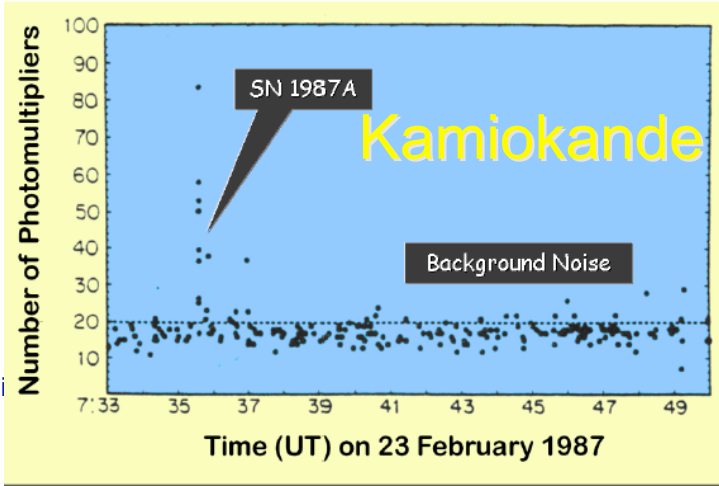
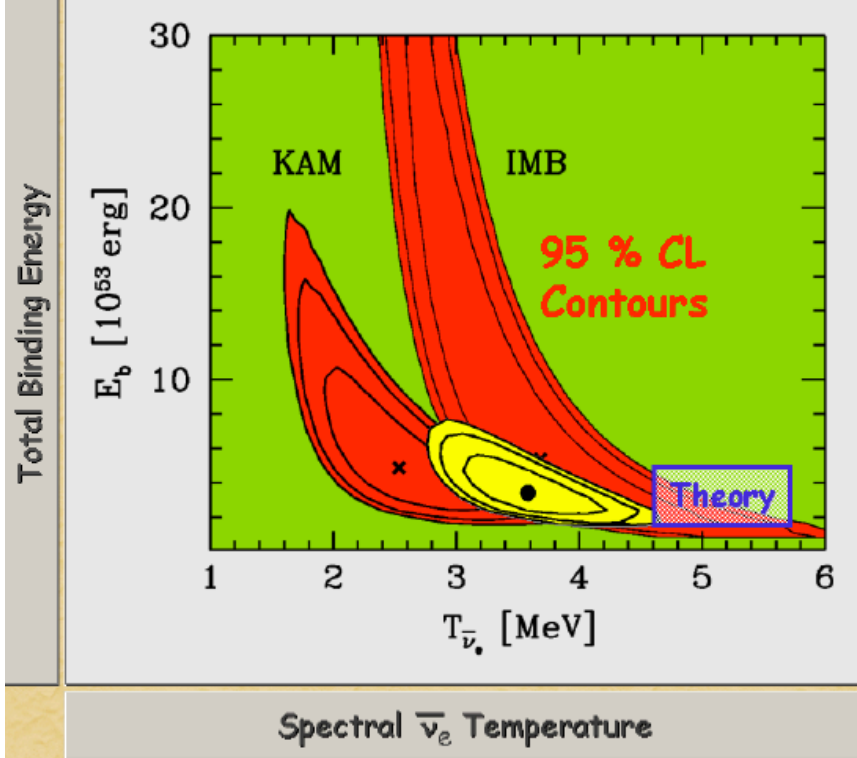
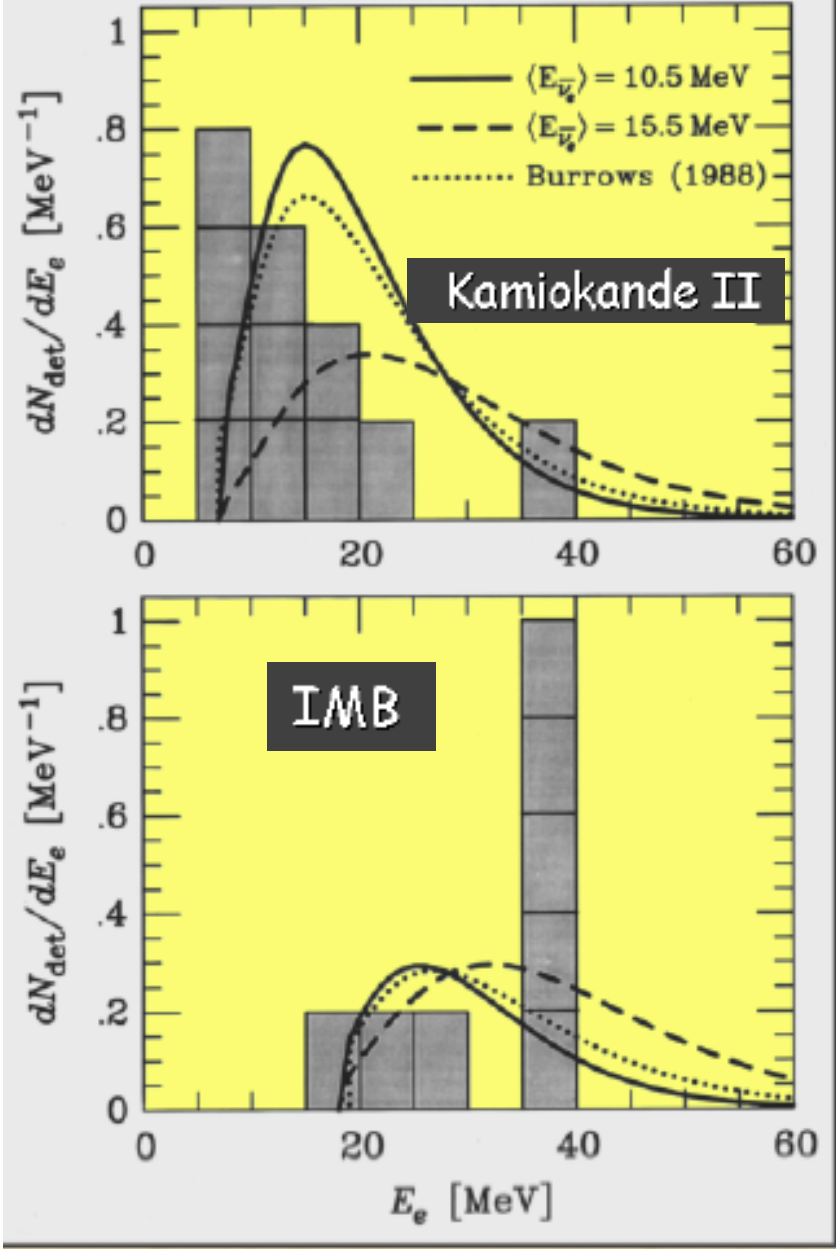
Scintillator detectors

LVD 1 kt (Jun 92 - Mar 03 - Jan 01 final configuration) 3511 d $E_{\text{th}} = 4\text{-}7$ MeV 0.2 SN/yr (90%c.l.) in Galaxy, in SNEWS since 98

Expected events from SN at 8.5 kpc 320 (210 in MACRO upper limit 0.27 SN/yr)

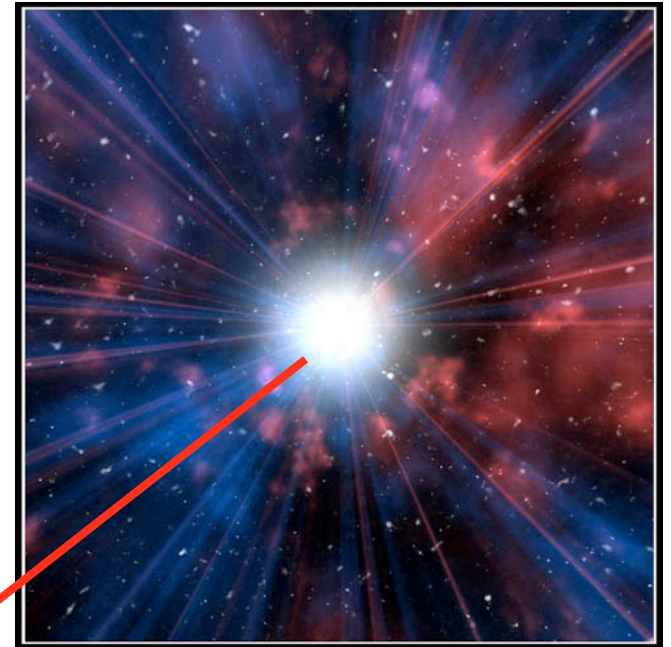
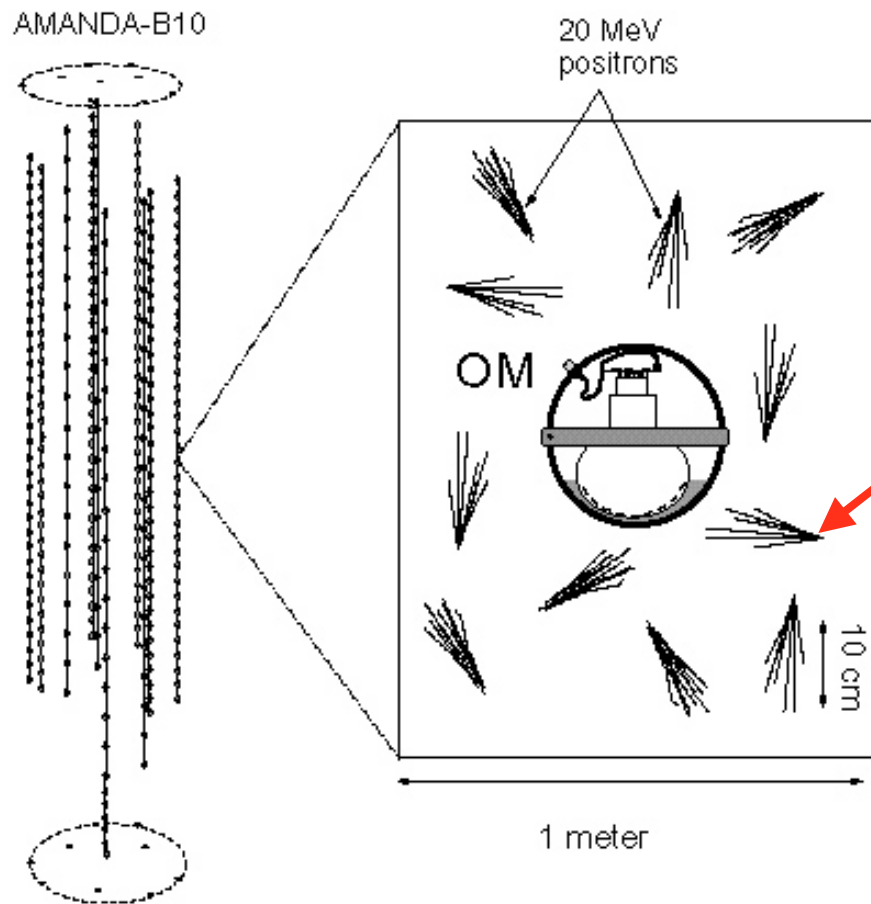
Others: Kamland (1 kt), MiniBoone (0.6 kt), Borexino (0.3 kt),...

SN1987A

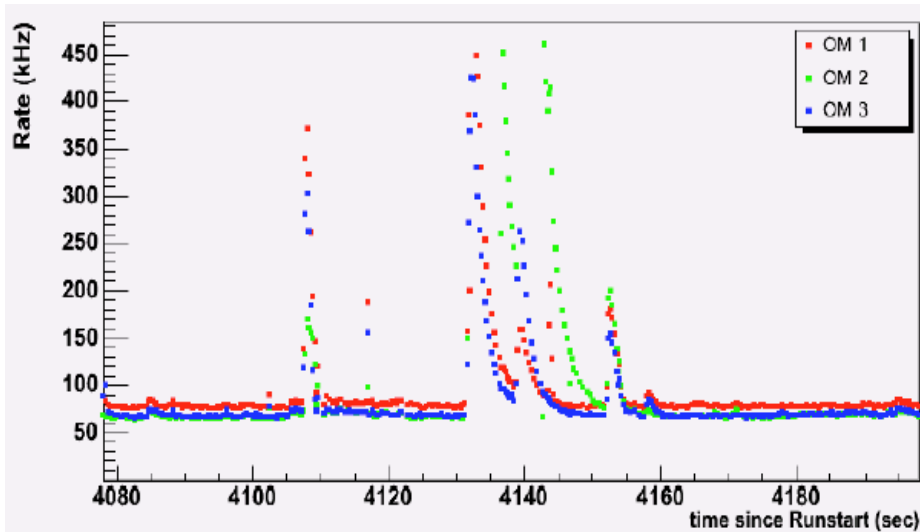


a Montaruli

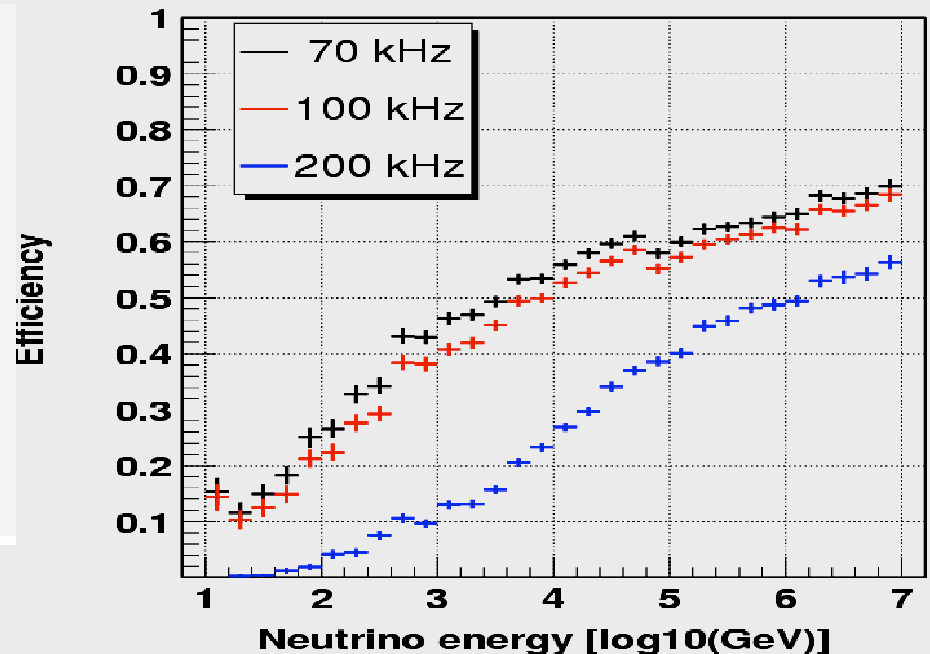
AMANDA as supernova detector



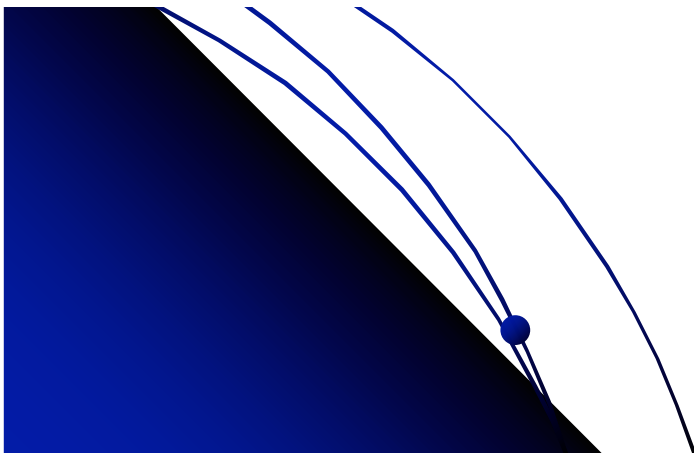
Optical Background and Filtering in ANTARES



Counting rate due to ^{40}K β decay and bacteria bioluminescence. Bursts are due to macro-organism



All data (>0.3 pe) sent to shore 1GB/s
Offline filter: 1 MB/s causality condition

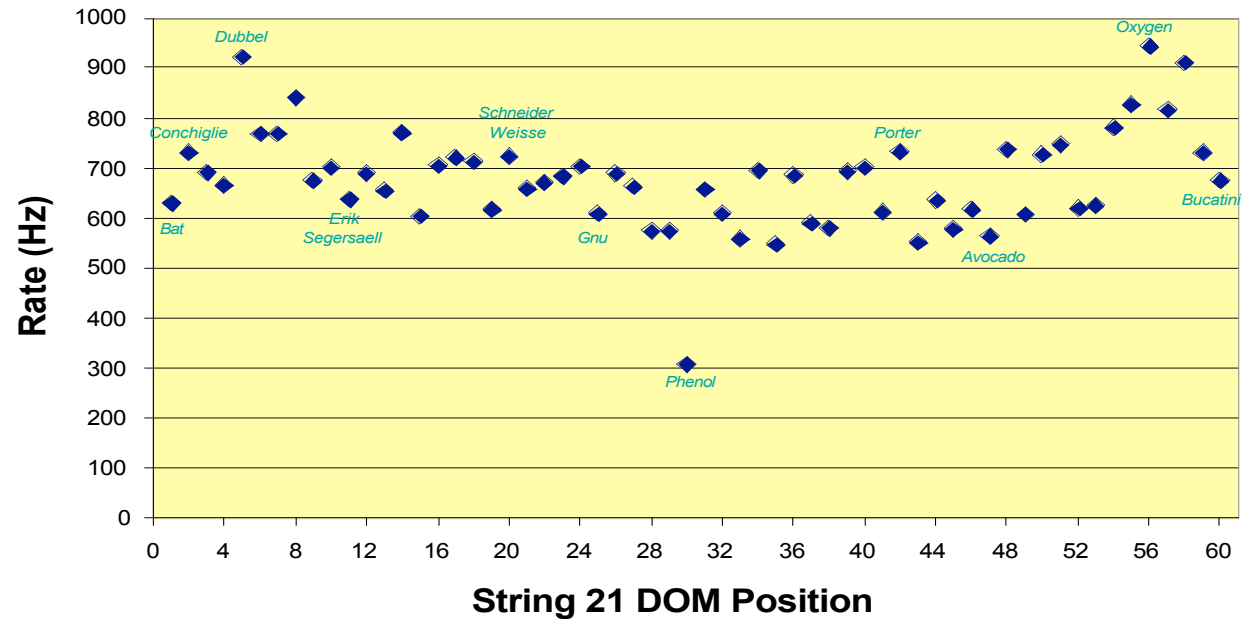


Ice is an extremely quite environment!

1st IceCube string

- The IceCube optical sensors were optimized for low noise.
- Research on glass material resulted in lower contamination with radioactivity.
- Fewer background photons from the glass.

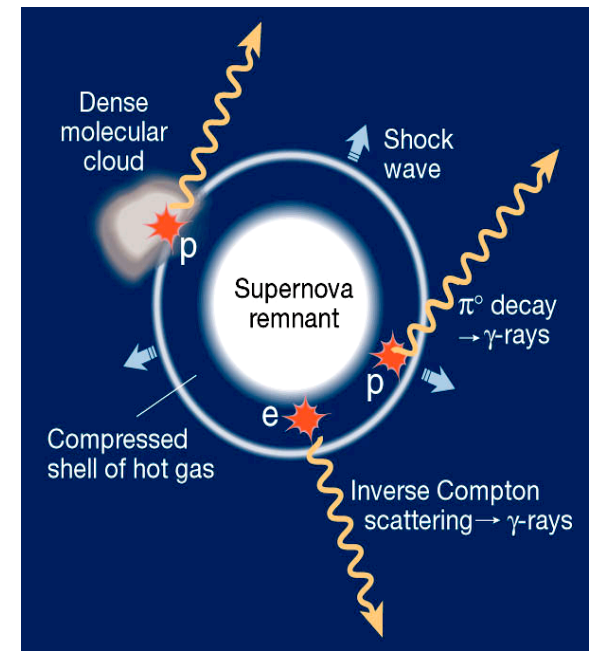
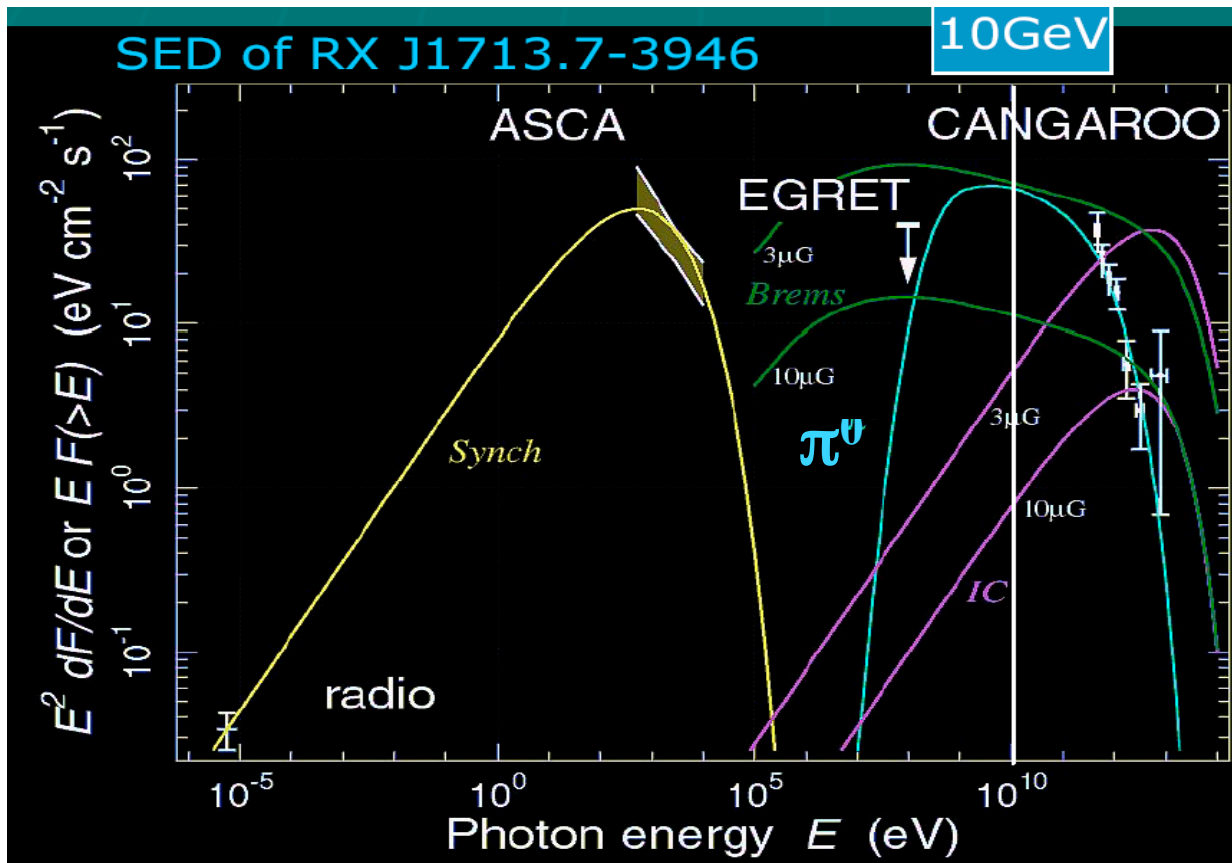
September 2005 InIce Noise Rates



Galactic point Sources

The case of RXJ1713.7-3946

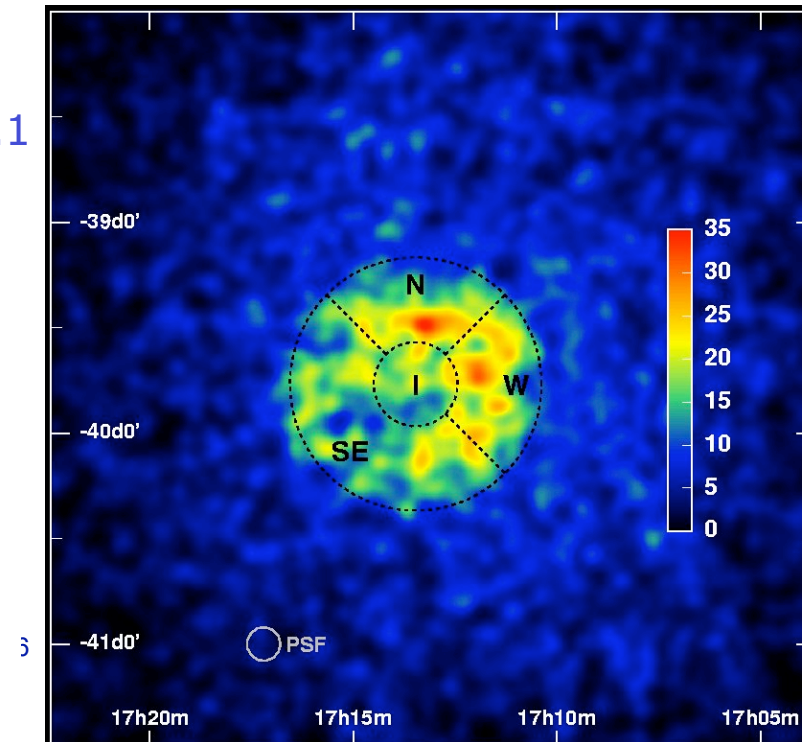
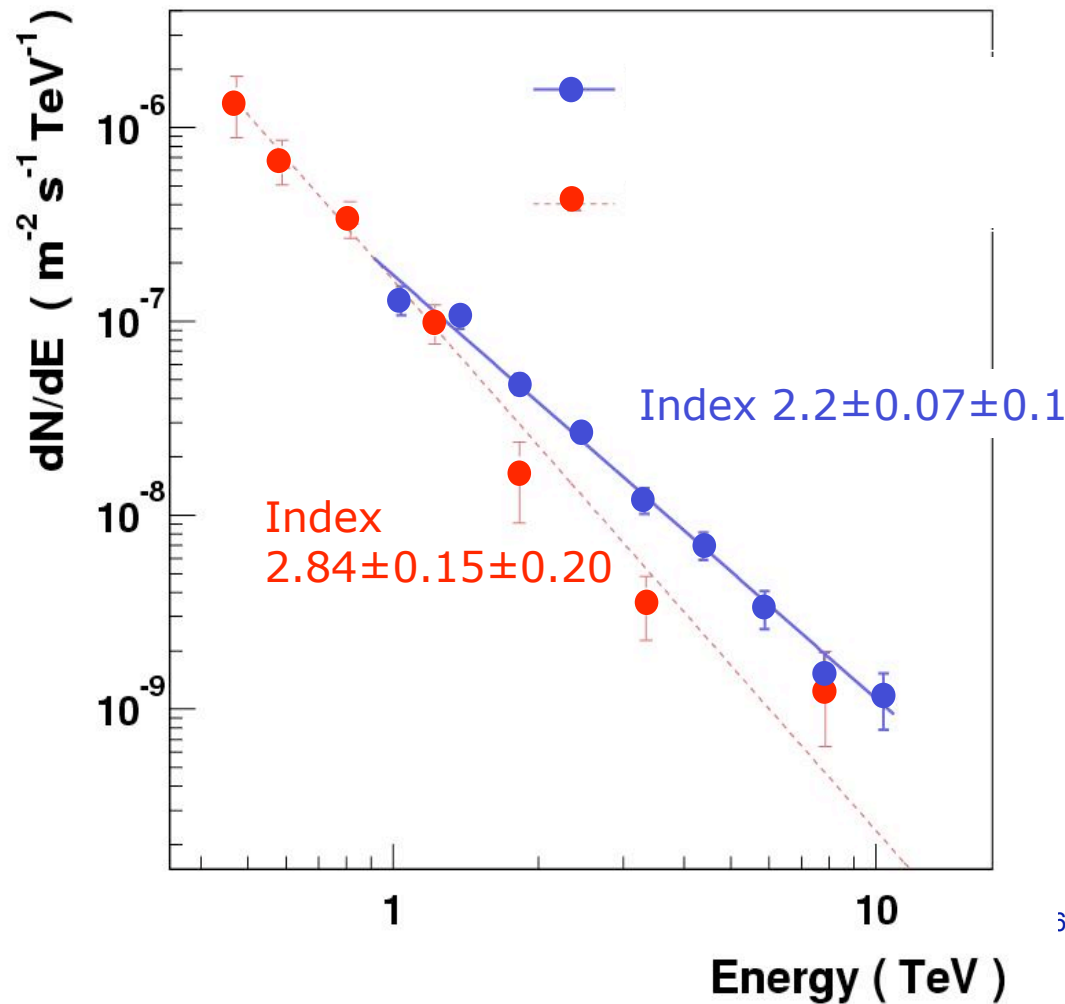
Open problem: elusive π^0 produced in accelerated nuclei collisions with SN ambient material Still not a clear evidence BUT...CANGAROO claim



Controversial
Reimer et al., A&A390,2002
Incompatible with EGRET

RXJ1713.7-3946

No cut-off in the HE tail of HESS spectrum favors π^0 decay scenario
respect to the case of em processes
Study of electron density and B can help



Predictions Galactic sources

Burgio, Bednarek, TM, New Astron. Rev. 49, 2005

Source Type	Distance (kpc)	E_ν (GeV)	N_{ν_e} ($\text{km}^{-2} \text{yr}^{-1}$)	Ref.
Supernovae Shocks pulsars	10	$\llsim 10^3$ $\sim 10^2 - 10^6$ $\sim 10^5 - 10^8$ $\sim 10 - 10^8$	~ 100 50 - 1000 $\sim 100 - 1000$ $\llsim 1000$	Waxman & Loeb 2001 Protheroe et al. 1998 Beall & Bednarek 2002 Nagataki 2004
Plerions Crab	0.5 - 4.4 2	$< 10^3 - 10^5$ $\sim 10^3 - 5 \cdot 10^5$ $\sim 10^3 - 5 \cdot 10^5$ $\sim 10^3 - 5 \cdot 10^5$ 10 - 10^6	$\sim 1 - 12$ $\llsim 1$ a few ~ 1 $\sim 4 - 14$	Guetta & Amatto 2003 Bednarek 2003 Bednarek & Protheroe 1997 Bednarek 2003 Amato et al. 2003
Shell SNRs SNR RX J1713.7 Sgr A East	6 8	$\llsim 10^4$ $\llsim 10^5$	~ 40 ~ 140	Alvarez-Muñiz & Halzen 2002
Pulsars + Clouds Galactic Centre Cygnus OB2	8 1.7	$10^4 - 10^7$ $> \sim 10^3$ $10^4 - 10^7$ $\llsim 10^6$	$\sim 2 - 30$ a few ~ 0.5 ~ 4	Bednarek 2002 Torres et al. 2004 Bednarek 2003 Anchordoqui et al. 2003
Binary systems A0535+26	2.6	$3 \cdot 10^2 - 10^3$	a few	Anchordoqui et al. 2003
Microquasars	1 - 10	$10^3 - 10^5$	1 - 300	Distefano et al. 2002
Magnetars	3 - 16	$\llsim 10^5$	1.7 ($0.1/\Delta\Omega$) ($5/d^2$)	Zhang et al. 2003

Microquasars

Neutrinos from p- γ interactions (photons from synchr. Emission of electrons accelerated in jet or from accretion disc) Distefano, Waxman et al 2002

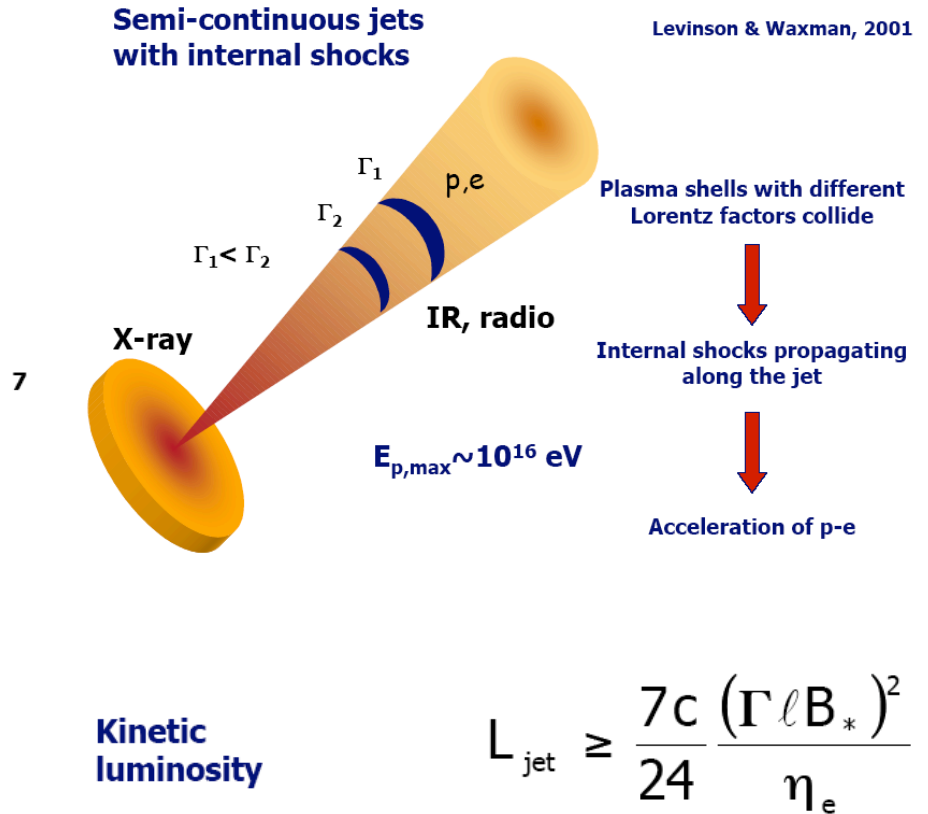
Neutrino flux at Earth

$$\begin{aligned}
 & \nu_{\mu} \text{ flux to the Earth} \\
 & f_{\nu_{\mu}} \sim \frac{1}{2} \eta_p f_{\pi} \delta^4 \frac{L_{\text{jet}} / 8}{4\pi D^2} \\
 & 1 \text{ TeV} \leq E_{\nu_{\mu}} \leq 100 \text{ TeV}
 \end{aligned}$$

- $\eta_p \sim 10\%$: fraction of the jet energy injected as Fermi protons
- L_{jet} : kinetic luminosity of the jet
- δ : jet Doppler factor $\delta = [\Gamma(1 - \beta \cos \theta)]^{-1}$
- D : source-Earth distance

$$\Gamma \sim 1$$

l : linear size of the blob

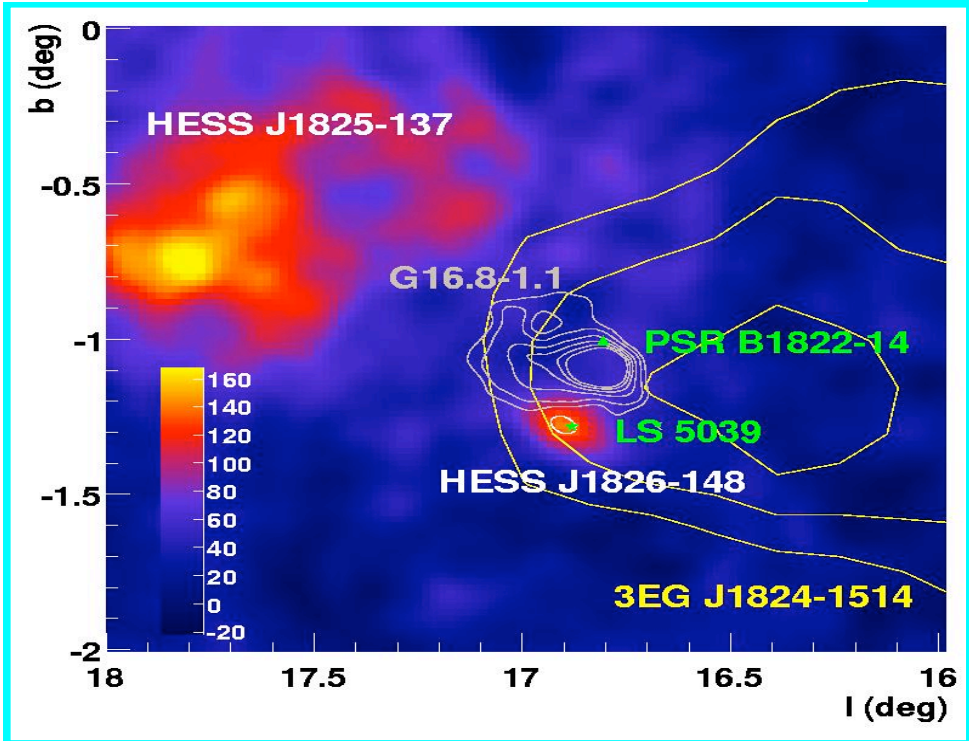
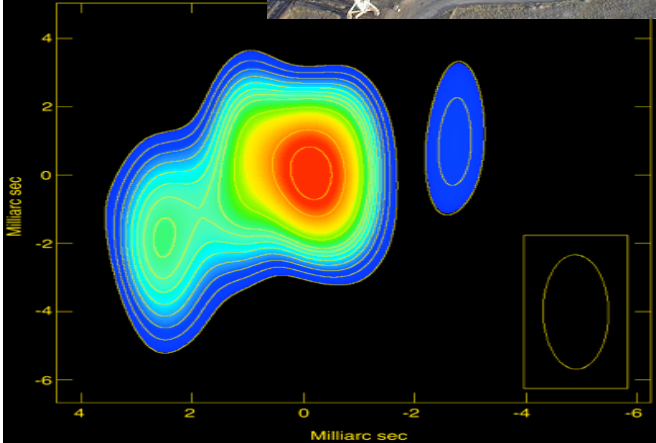
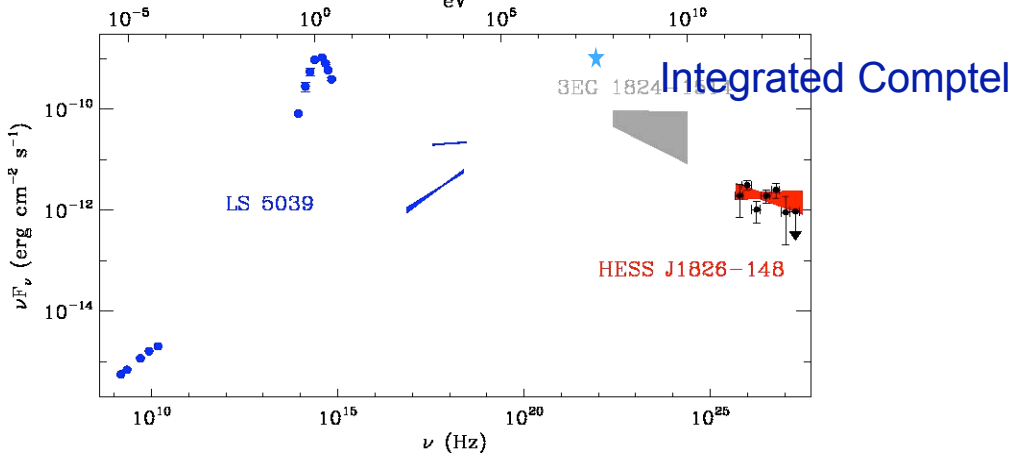


7



LS5039

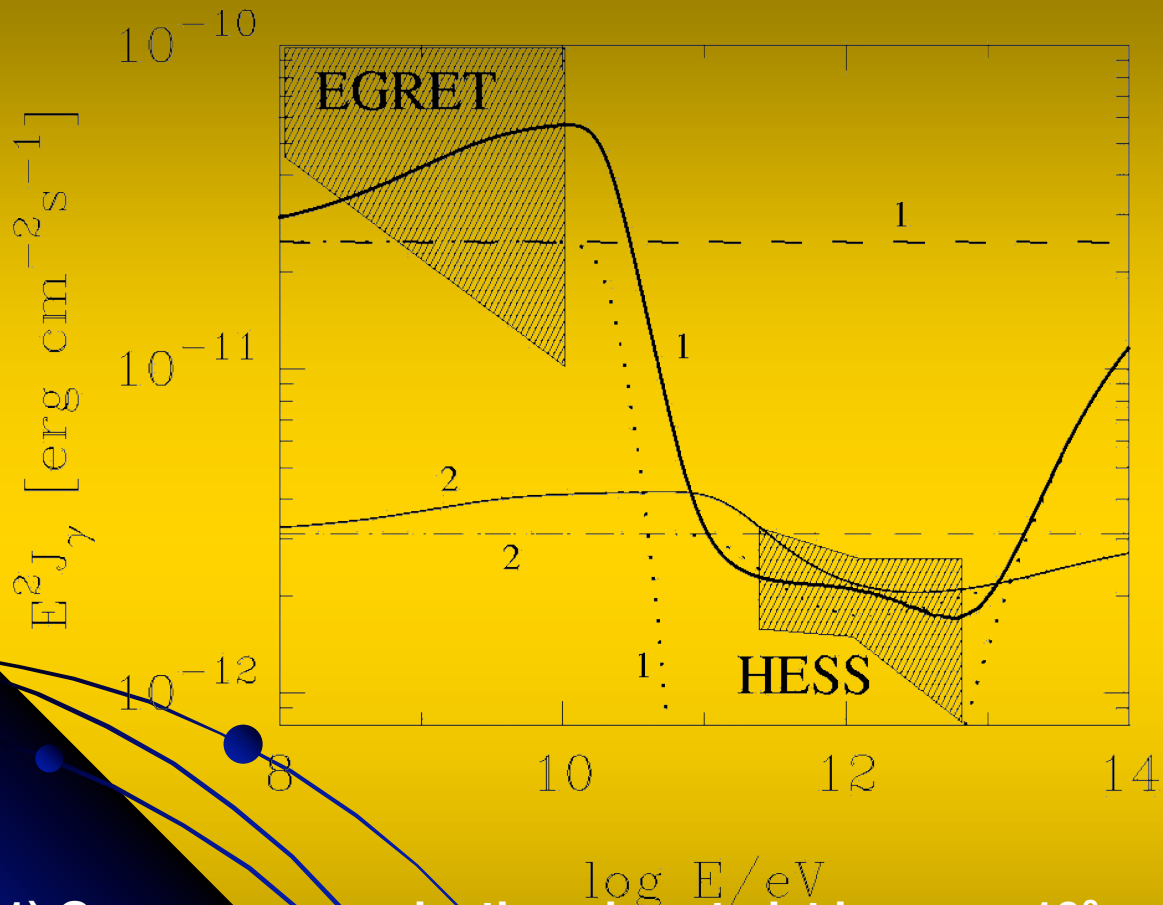
Very Large Array



- Source angular size ~ 50 arcsec
- Source distance ~ 2.5 kpc
- Gamma rays within radius ~ 0.6 pc
- Likely to be associated to
 - 3EG J1824-1514
 - Hard E⁻² spectrum

Apr. 2006

Time averaged gamma ray spectra

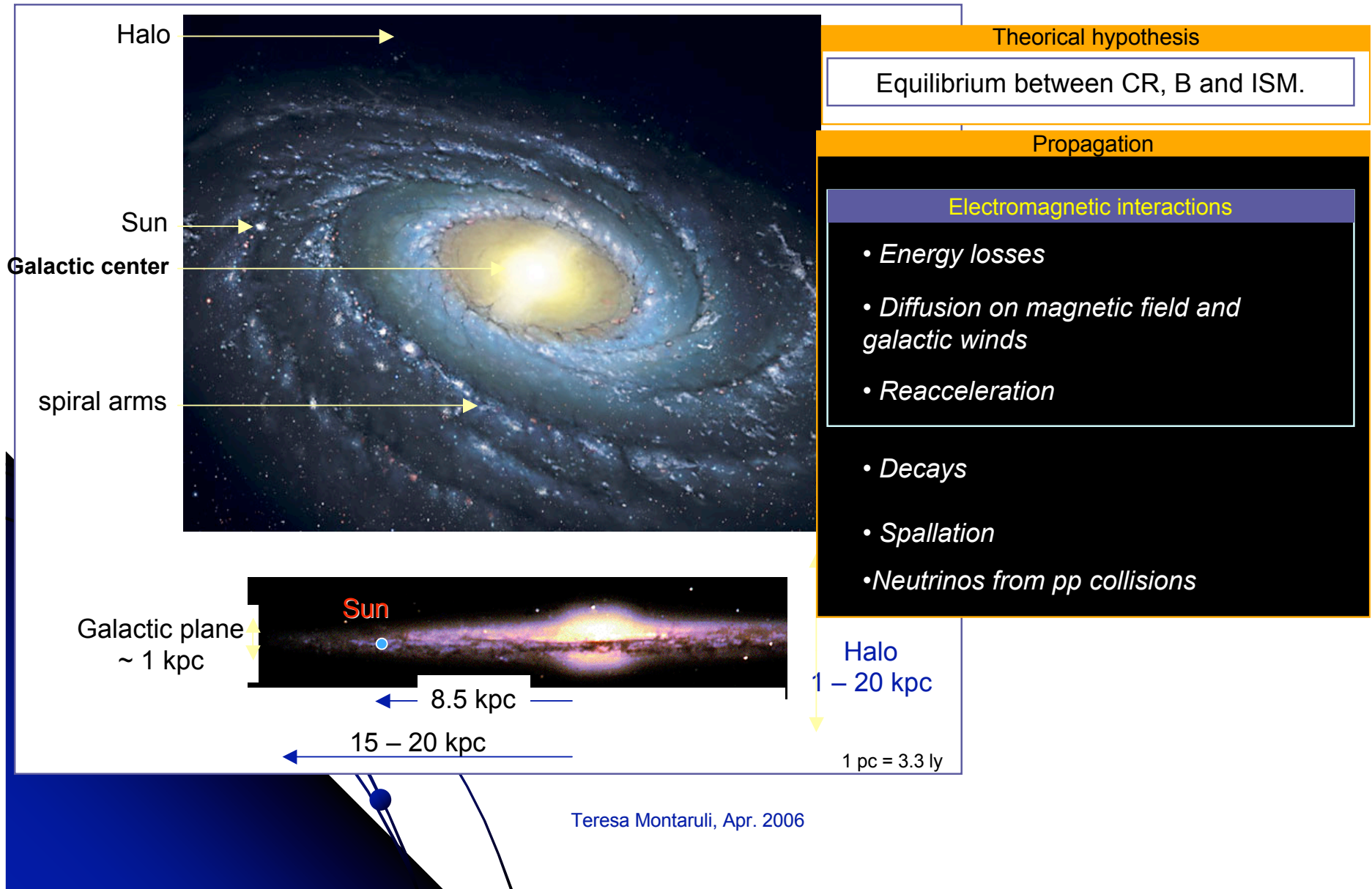


Absorption by
synchrotron
cascading ____

(1) Gamma ray production close to jet base $z < 10^8$ cm (more absorption, larger flux at low energy)

(2) Gamma ray production far from jet base $z < 10^{13}$ cm

The Galaxy



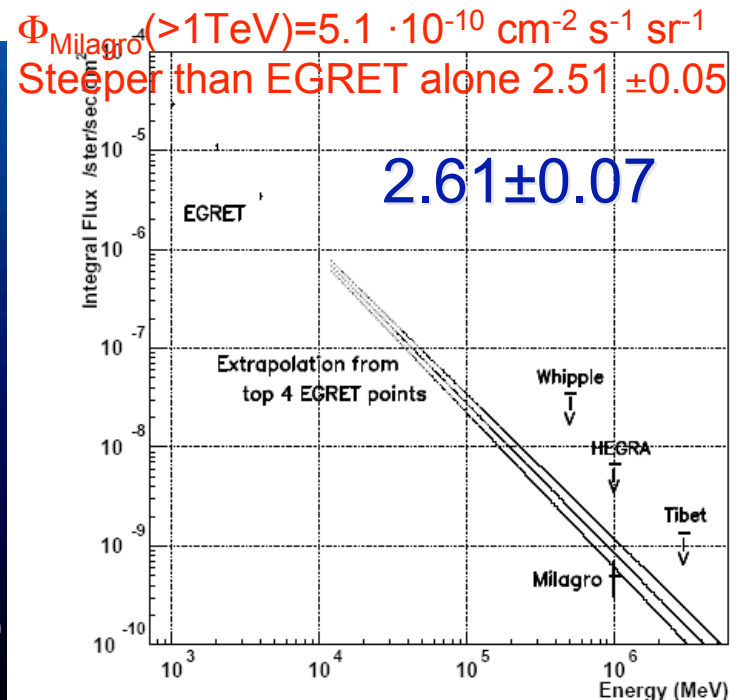
γ observations

- EGRET observed a diffuse emission 100MeV-10 GeV from Galactic Centre region (300 pc): excess > factor 10 around 1 GeV
- *INTEGRAL*: resolved 91 point sources. 90% of 'diffuse' flux can be due to point sources <100 keV
- Milagro: discovery of TeV emission (astro-ph/0502303)
4.5 σ excess from $|b| < 5^\circ$ and $l \in [40^\circ, 100^\circ]$

Covered pond with 2 layers of PMTs, from relative timing 0.75° shower direction resolution, gamma-hadron discrimination based on shape of Cherenkov light emitted by showers

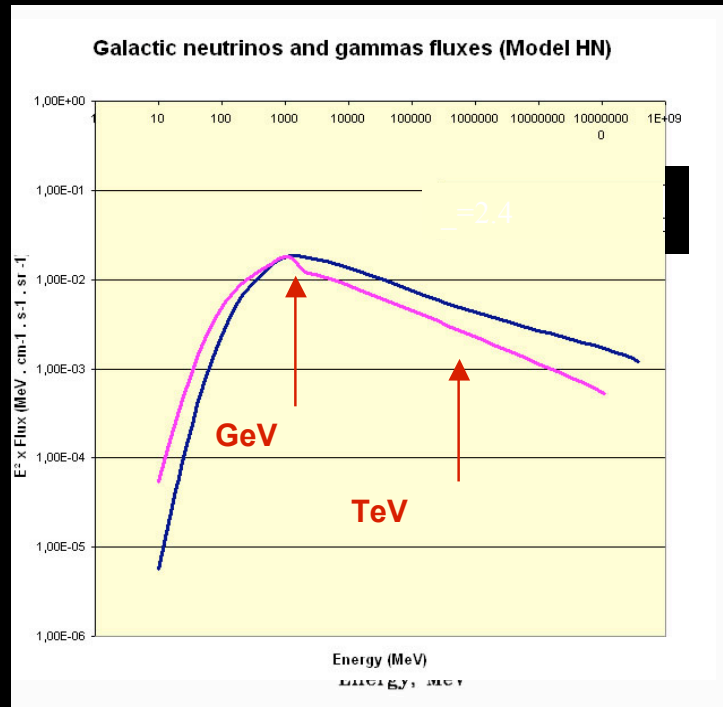


Lebrun et al. 2004 Nature 428, 293

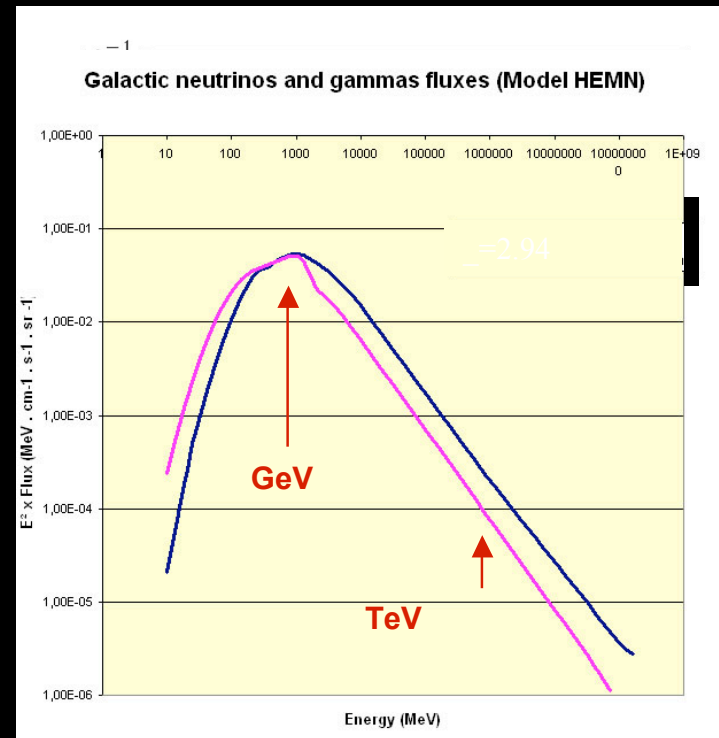


Extreme Models

Hard nucleus model $E^{-2.4}$



Hard electron model $E^{-2.9}$



For $E^{-2.4}$ 20 years of ANTARES to have 88% discovery prob

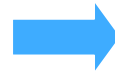
$$\frac{d\phi(E)}{dE} = A.E^{-\gamma}$$

— Gamma from π^0

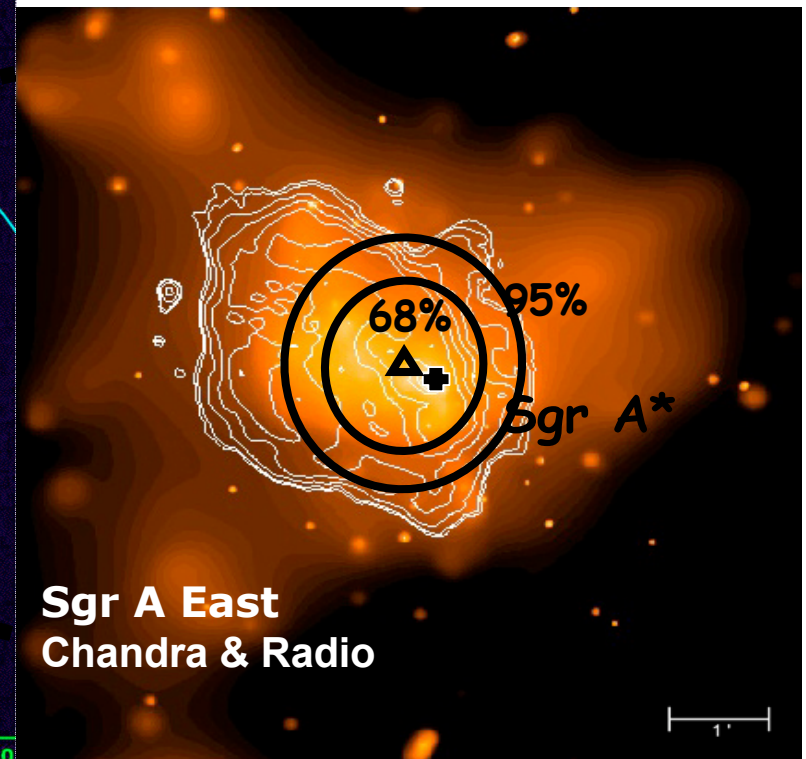
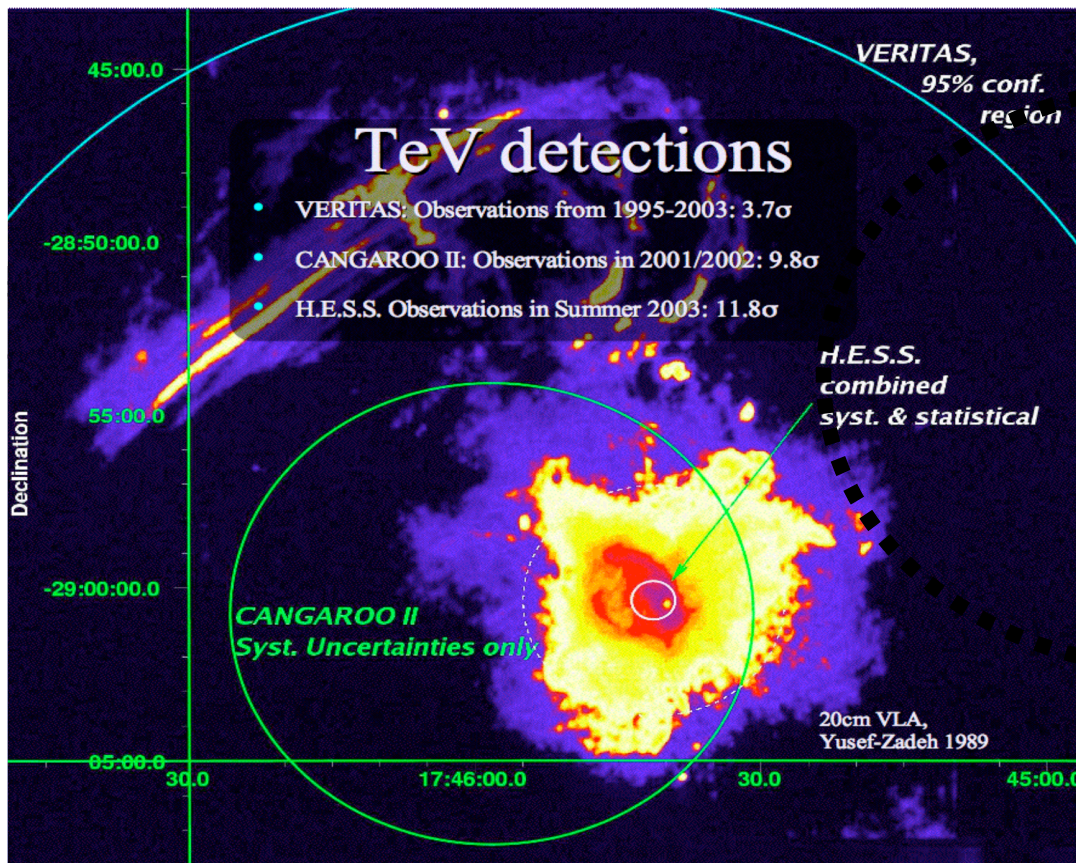
— Nu mu + anti nu mu

Galactic Centre

- High matter density and activity
- compact radio source Sgr A* possibly associated to black hole $\sim 3 \cdot 10^6 M_{\text{sun}}$ in the center
- Sgr A East SNR



HESS TeV- γ spectrum in disagreement with the other experiments Variability? localization? HESS 1 arcmin

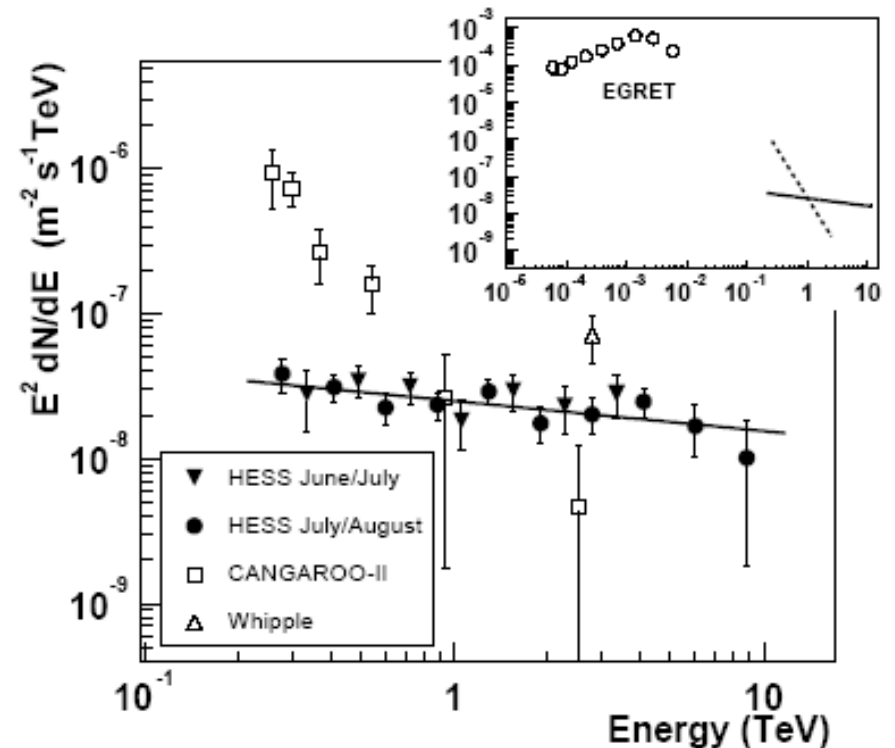


High Energy Stereoscopic System

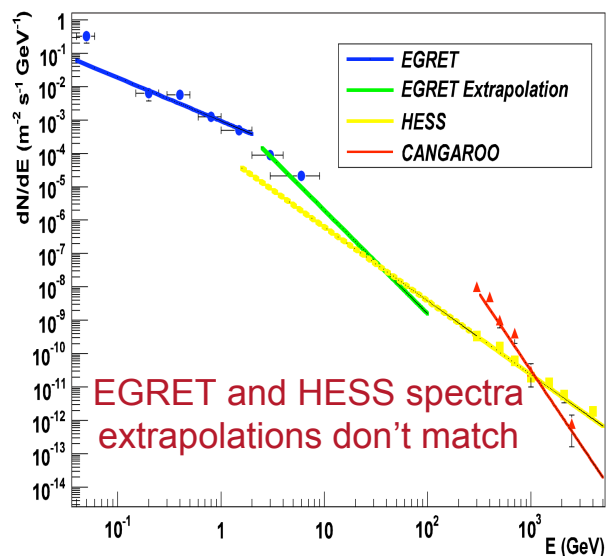
Four 12 m diameter telescopes running since ~ 1yr in Namibia (+1 large) $E_{th} \sim 100$ GeV

Cherenkov light is emitted by showers induced by high-energy gamma rays This light is very faint - about $10 \gamma/s/m^2$ at $E_\gamma=100$ GeV - and the duration of the light flash is only a few nsec. Large mirrors, fast photon detectors and short signal-integration times are required to collect enough light from the shower, with minimal contamination from night-sky background light.

γ direction $< 0.1^\circ$



Neutrinos flux with different constraints



$E_\mu > 1 \text{ TeV}$
 $d\Omega < 0.5^\circ$

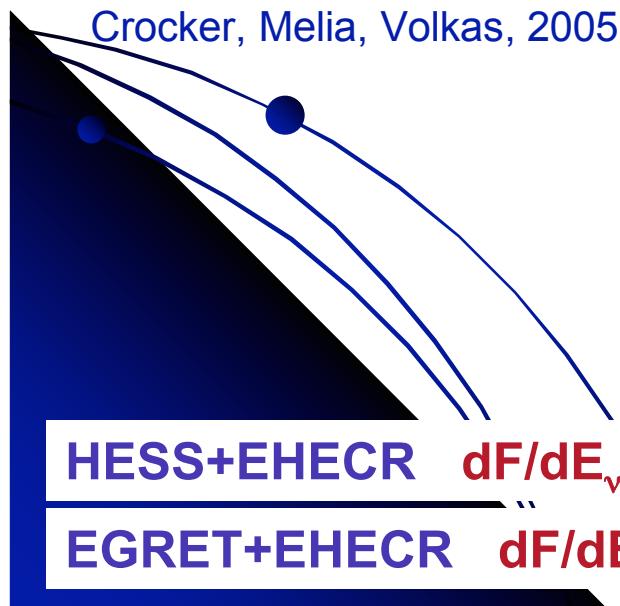
Not optimized

PRELIMINARY

ANTARES

KM3

		Signal events /year	Bkg events /year	Time for detection (4 σ CL)	Time for detection (4 σ CL)
GC	HESS	$2 \cdot 10^{-2}$		247 yr	6.2 yr
	HESS+EHECR	$5 \cdot 10^{-2}$	$7 \cdot 10^{-3}$	64 yr	1.7 yr
	EGRET+EHECR	2.6		0.4 yr	week



RX J1713.7 -3946	Constantini et al (2005)	0.16		15 yr	0.4 yr
	Halzen et al (2002)	0.95	$9 \cdot 10^{-3}$	1.8 yr	3 weeks
PSR B1509- 58	HESS	0.11	$1.1 \cdot 10^{-2}$	22 yr	0.6 yr

HESS+EHECR $dF/dE_{\nu\mu+\bar{\nu}\mu} = 1.3 \times 10^{-5} E^{-2.0} \text{ m}^{-2} \text{ s}^{-1} \text{ GeV}^{-1}$

EGRET+EHECR $dF/dE_{\nu\mu+\bar{\nu}\mu} = 4.1 \times 10^{-3} E^{-2.22} \text{ m}^{-2} \text{ s}^{-1} \text{ GeV}^{-1}$

Upper bounds on X-galactic fluxes

Cosmic p accelerators produce CRs, γ 's and ν 's

Ultimate bound of any scenario involving ν and γ production from π s: diffuse extra-galactic γ background $E^2 F_\nu < 6 \cdot 10^{-7} \text{ GeV} / \text{cm}^2 \text{ s sr}$ (EGRET)

Measured UHECR flux provides most restrictive limit (Waxman & Bahcall (1999))

- optically thin sources: nucleons from photohadronic interactions escape
- CR flux above the ankle ($> 3 \cdot 10^{18} \text{ eV}$) are extragalactic protons with E^{-2} spectrum $\Rightarrow E^2 F_\nu < 4.5 \cdot 10^{-8} \text{ GeV} / (\text{cm}^2 \text{ s sr})$

This bound does not apply to harder spectra or optically thick

Mannheim, Protheroe & Rachen (2000):
Magnetic fields and uncertainties in photohadronic interactions of protons can largely affect the bound as these effects restrict number of protons able to escape

