

How do we interpret the gravitational redshift?

Observers at emitter and receiver measure time differently.

Gravitational time dilation

$$\omega \sim \frac{1}{d\tau}$$

1  $d\tau_1$

$$\frac{d\tau_1 - d\tau_2}{d\tau_2} = \frac{\phi_1 - \phi_2}{c^2}$$

2  $d\tau_2$

$$\Rightarrow d\tau_1 = \left( 1 + \frac{\phi_1 - \phi_2}{c^2} \right) d\tau_2$$

$$= \left( 1 + \frac{\Delta\phi}{c^2} \right) d\tau_2$$

time dilation

For static gravitational field, can integrate to get

$$\tau_1 = \left( 1 + \frac{\phi_1 - \phi_2}{c^2} \right) \tau_2$$

Higher potential runs faster

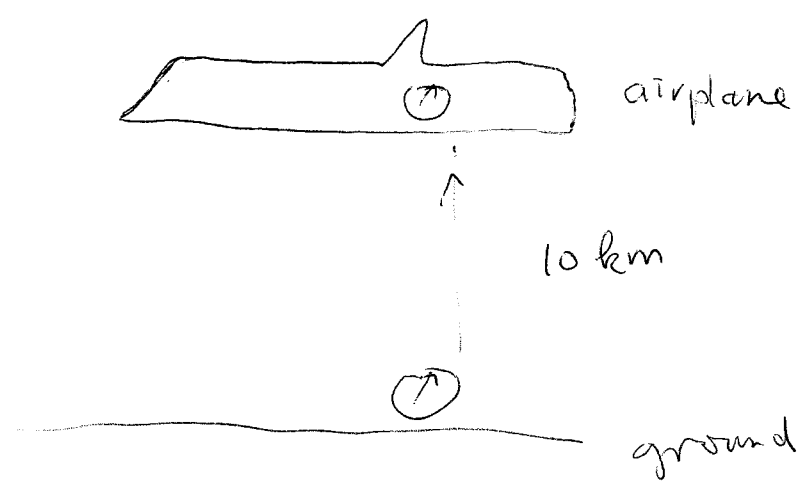
Note : ① Time dilation even though there is no relative motion

② Unlike time dilation in SR (ie. each see other's clock run slower)

Here :

lower clock see upper clock run fast  
upper clock " lower " " " slow

③ Time dilation by gravity has been tested

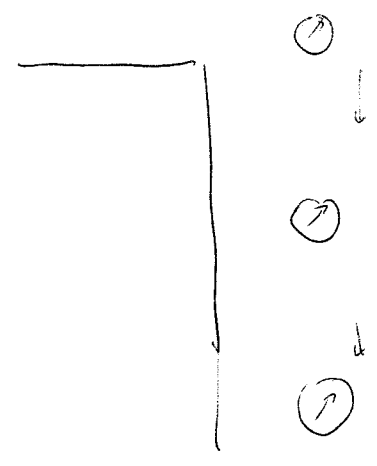


After correcting for SR effect, clock in airplane still runs fast

Also: applications in GPS (homework)

# More direct derivation of time dilation

Drop a clock from a building



Go to free-falling frame (inertial frame)

⇒ SR applies

$$t_1^{ff} = \gamma_1 \tau_1 \quad t_2^{ff} = \gamma_2 \tau_2$$

↑ free fall frame      ↑ rest frame

We are interested in comparing  $d\tau_1$  and  $d\tau_2$  where  $dt_1^{ff} = dt_2^{ff}$

$$\frac{d\tau_1}{d\tau_2} = \sqrt{\frac{1 - u_1^2/c^2}{1 - u_2^2/c^2}} \approx 1 - \frac{1}{2} \left( \frac{u_1^2 - u_2^2}{c^2} \right) + \dots$$

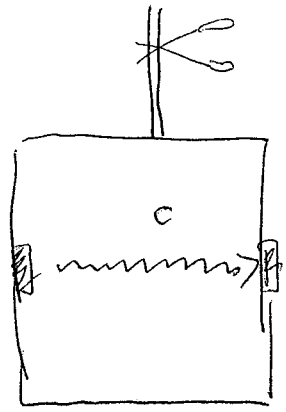
$$= 1 + \frac{\phi_1 - \phi_2}{c^2}$$

Conservation of energy

$$\frac{1}{2} m u_i^2 + \phi_i = \text{constant}$$

# Bending of light

First, qualitative:



Free-falling frame = inertial frame

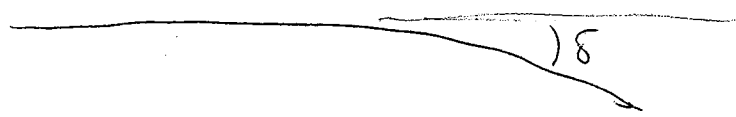
light travels horizontally with speed =  $c$

Now according to our frame, elevator is dropping so

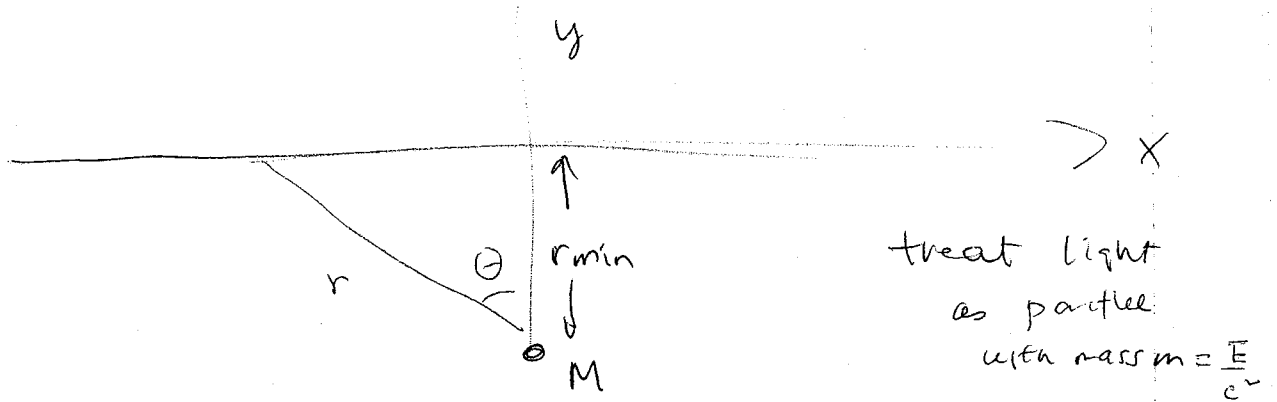


light bends like a parabola under influence of gravity

We want to calculate the angle  $\delta$  light is bent in the presence of a mass  $M$



Let us get the answer first by cheating



$$\Delta P_y = \int F_y dt = \int F \cos \theta dt$$

$$= \int \frac{GMm}{r^2} \cos \theta dt$$

$$= \int \frac{GMm}{r^2} (\cos \theta) \frac{dx}{c}$$

$$= \int_{-\infty}^{\infty} \frac{GMm}{c} \frac{\cos \theta dx}{x^2 + r_{\min}^2}$$

$$= \frac{GMm}{c} \int_{-\infty}^{\infty} \frac{\cos \theta dx}{x^2 + r_{\min}^2}$$

$$= \frac{GMm}{r_{\min} c} \int_{-\pi/2}^{\pi/2} \cos \theta d\theta$$

$$x = r_{\min} \tan \theta$$

$$dx = r_{\min} \sec^2 \theta d\theta$$

$$= \frac{2GMm}{r_{\min} c}$$

$$\frac{\Delta P_y}{P} = \frac{2GMm}{r_{min}c} \frac{1}{E/c} = \frac{2GM}{r_{min}c^2}$$

$\frac{E}{c^2}$

Other than the fact that this is not the full story for GR, we got this result just because we were lucky.

Problems with this method

- ① Massless particle only makes sense in relativity, do not expect  $\frac{GMm}{r^2}$  to work
- ② Derivation does not take into account gravitational time dilation, which is crucial in the reasoning behind bending of light.

## Bending of light (from EP)

Speed of light  $\frac{dr}{d\tau} = c$  universal constant

According to an observer at a different position

speed appears to be different

because gravitational potential is different.

Let's choose the observer at infinity

$$t = \tau(\infty) \quad \Phi(\infty) = 0$$

$$d\tau = \left(1 + \frac{\Phi(r)}{c^2}\right) dt$$

Apparent speed of light

$$\begin{aligned} c(r) &\equiv \frac{dr}{dt} = \left(1 + \frac{\Phi(r)}{c^2}\right) \frac{dr}{d\tau} = \left(1 + \frac{\Phi(r)}{c^2}\right) c \\ &\equiv \frac{c}{n(r)} \end{aligned}$$

where  $n(r)$  can be interpreted as an index of refraction

$$n(r) \equiv \left(1 + \frac{\Phi(r)}{c^2}\right)^{-1} \approx 1 - \frac{\Phi(r)}{c^2}$$

## Important to note

- $c = \frac{dr}{dt}$  is still a universal constant

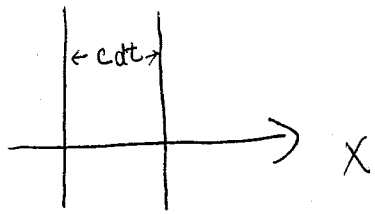
rather it is the clock which runs differently  
so it appears to change according to a  
fixed clock

- A dramatic example: Blackhole

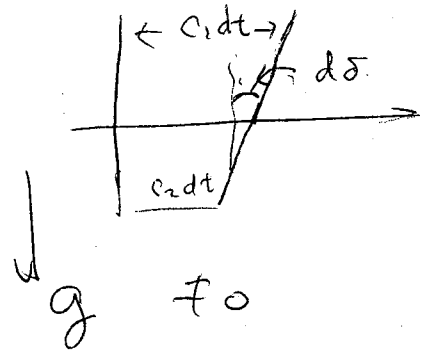
Infinite time dilation  $\Rightarrow$  appears to an  
observer at  $\infty$  that  
it takes  $\infty$  amount  
of time to leave the  
black hole even  
though proper time  
is finite.



## Bending of light (the EP expectation)



$$g = 0$$



$$g \neq 0$$

$$d\delta \approx \frac{(c_1 - c_2) dt}{dy} \approx \frac{dc(r)}{dy} \frac{dx}{c}$$

$$= \left( \frac{\partial n^{-1}}{\partial y} \right) dx$$

$$= - \frac{1}{n^2} \left( \frac{\partial n}{\partial y} \right) dx$$

$$\approx - \left( \frac{\partial n}{\partial y} \right) dx$$

$$n \approx 1 - \frac{\phi(r)}{c^2}$$

$$\frac{\partial n}{\partial y} = \frac{\partial n}{\partial r} \frac{\partial r}{\partial y} = - \frac{1}{c^2} \left( \frac{\partial \phi}{\partial r} \right) \frac{y}{r}$$

$$= - \frac{GM}{c^2 r^2} \frac{y}{r}$$

$$\delta = \int_{-\infty}^{\infty} \frac{GM y}{c^2 r^3} dx = \frac{GM}{c^2} \int_{-\infty}^{\infty} \frac{y dx}{r^3}$$

For small bending,  $y \approx r_{\min}$

$$\Rightarrow \delta = \frac{GM}{c^2} \int \frac{r_{\min} dx}{(x^2 + r_{\min}^2)^{3/2}}$$

$$= \frac{2GM}{c^2 r_{\min}}$$

For  $r_{\min} = R_{\odot}$        $M = M_{\odot}$   
 $\uparrow$   
 sun

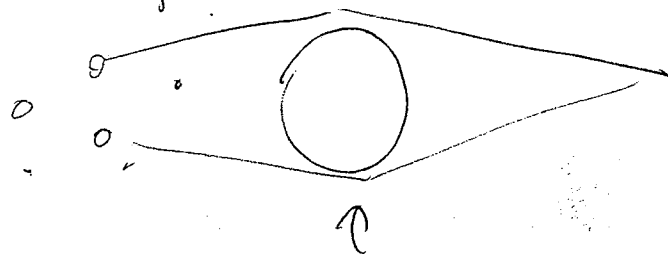
$\Rightarrow \delta = 0.875''$       turns out only  $1/2$   
 of the correct answer  
 for GR

Experimental verification of such (correct)  
 angle of deflection is one of the greatest  
 triumphs of GR

### Application

Gravitational  
 lensing

galaxy cluster



$\uparrow$   
 massive object acts  
 as a lens

see distinct  
 galaxies  
 that  
 are  
 otherwise  
 out of  
 sight