

Particle in a box :  $E_n = \frac{\pi^2 \hbar^2}{2mL^2} n^2 \quad n=1, 2, 3, \dots$

Harmonic oscillator  $E_n = (n + \frac{1}{2}) \hbar \omega \quad n=0, 1, 2, \dots$

More generally,

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(x)\psi = E\psi$$

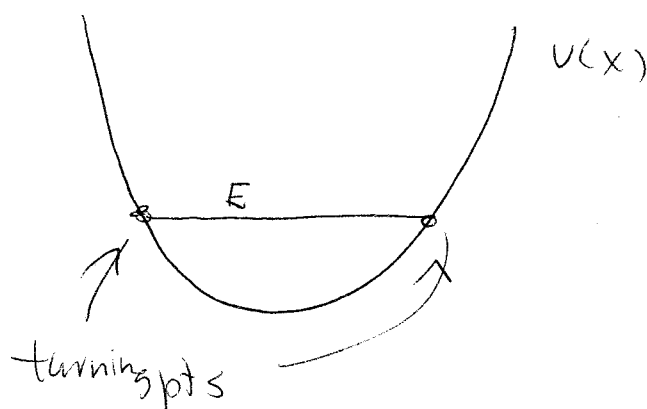
$$\frac{d^2\psi}{dx^2} = -\frac{2m}{\hbar^2} (E - V)\psi$$

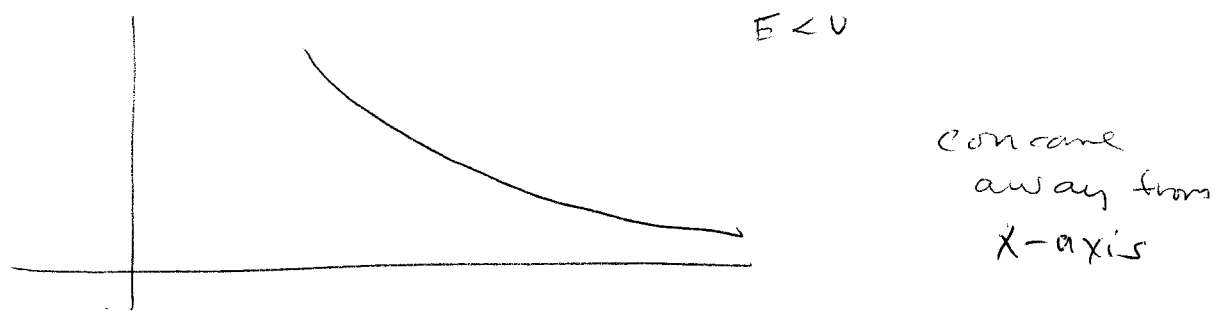
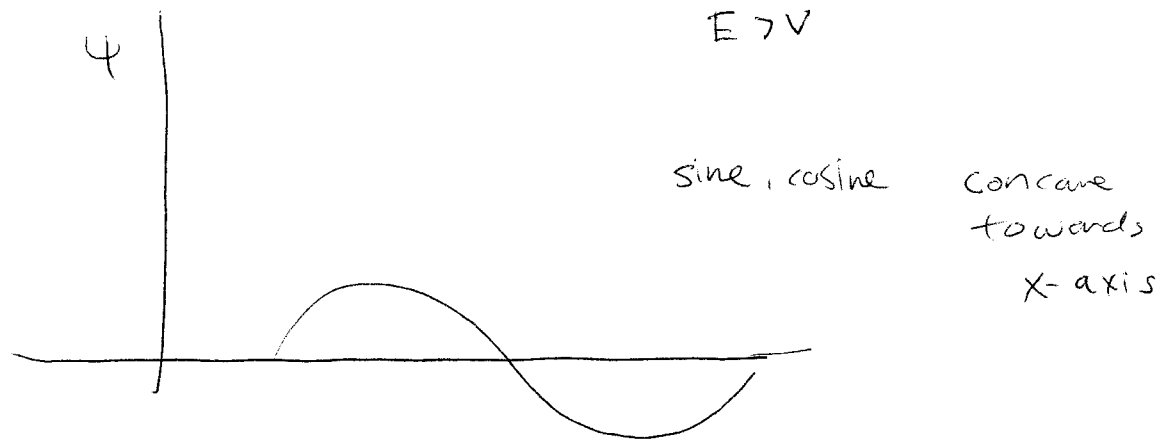
If  $E > V(x)$   $\frac{d^2\psi}{dx^2} = -k^2(x)\psi$  where  $k^2(x) = \frac{2m}{\hbar^2} (E - V)$

$\Rightarrow \psi$  looks like sine & cosine

$E < V(x)$   $\frac{d^2\psi}{dx^2} = \alpha^2(x)\psi$  where  $\alpha^2(x) = \frac{2m}{\hbar^2} (V - E)$

$$\Rightarrow \psi = e^{\pm \alpha x}$$



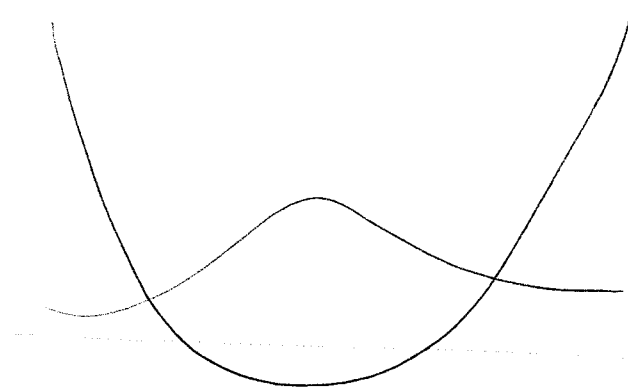


These qualitative features together with

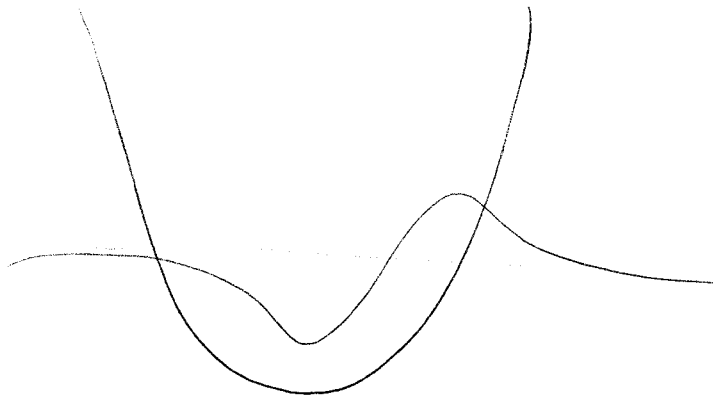
- ① parity of wavefunction
- ②  $n$ -th excited state has  $n$  zero crossings
- ③  $P(x) \propto \frac{1}{v}$  classically

we can get a rather good idea of  $\psi(x)$

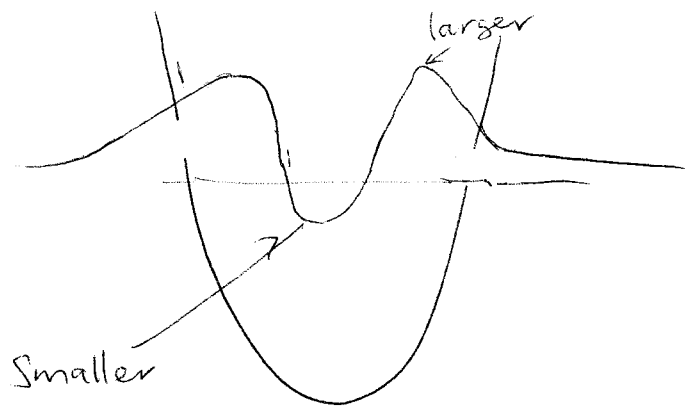
$n=0$



$n=1$



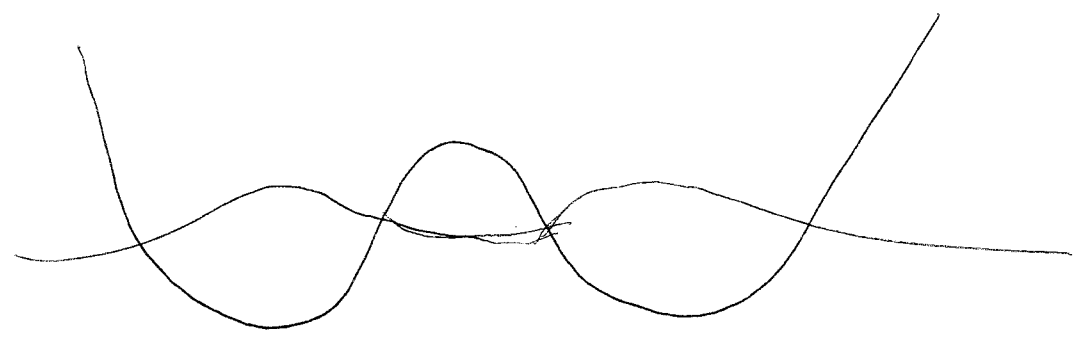
$n=2$



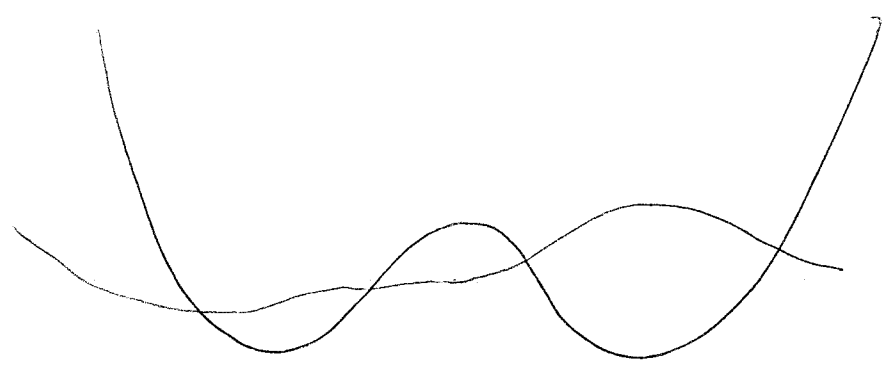
wavefunction should have smaller magnitude at the middle where the classical velocity is large.

# Double well potential

$n=0$



$n=1$



(5)

## Other simple generalizations

3D SHO :

$$V(x) = \frac{1}{2} k (x^2 + y^2 + z^2) = \frac{1}{2} m\omega^2 (x^2 + y^2 + z^2)$$

$$E = (n_1 + \frac{1}{2}) \hbar\omega + (n_2 + \frac{1}{2}) \hbar\omega + (n_3 + \frac{1}{2}) \hbar\omega$$

Energy levels are degenerate

Ground state  $n_1 = n_2 = n_3 = 0$   $E_0 = \frac{3}{2} \hbar\omega$  non-degenerate

but

First excited state  $n_1 = 1, n_2 = 0, n_3 = 0$   
(or permutations)

$$E_1 = \frac{5}{2} \hbar\omega \quad \underline{\underline{3\text{-fold degenerate}}}$$

More degeneracy at large  $E$ .